# Custodial Leptons and Higgs Decays

### Florian Goertz



Eidgenössische Technische Hochschule Zürich Swiss Federal Institute of Technology Zurich

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Carmona, FG JHEP 04(2013)163

### Composite Higgs Models

• Higgs is composite at small distances  $E\gg 1/I_H\sim f$  Kaplan, Georgi, Dimopoulos,...

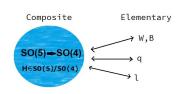


 $\Rightarrow$   $m_H$  saturated in IR  $\Rightarrow$  Hierarchy Problem solved

ullet Higgs as (pseudo-)Goldstone Boson  $\Rightarrow$   $m_H \ll m_
ho$  [like QCD pions]

 Minimal viable symmetry-breaking pattern:

⇒ Custodial Symmetry
The Minimal Composite Higgs Model (MCHM),
Agashe, Contino, Pomarol

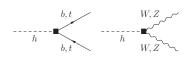


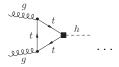
## Higgs Production and Decay in Composite Models

$$\mathsf{Higgs} = \mathsf{Goldstone} \to \mathsf{non-linear} \ \mathsf{realization} \ \Sigma_{\mathit{I}} = \left( \mathsf{Exp}[iH_{\hat{\mathsf{a}}}T^{\hat{\mathsf{a}}}/f] \right)_{\mathit{I5}}$$

Higgs decay constant  $f = m_
ho/g_
ho < m_
ho$ 

 $\Rightarrow$  Modification of Higgs couplings due to Goldstone nature, scale as trigonometric functions of  $v/f \Rightarrow indirect$  signs of model





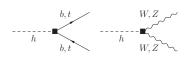
$$\kappa_{p}\equivrac{g_{hpp}}{g_{hpp}^{
m SM}}$$

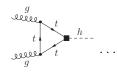
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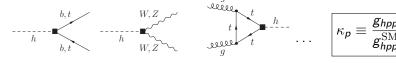
ATLAS Preliminary	m <sub>H</sub> = 125.5 GeV
W,Z H → bb W-7 lox [ca - 47 lo* Vi - 6 lot [ca - 47 lo*	_
H → cr vi=1 fire (cr=4.6 m² vi=6 m² (cr=10 m²	<b>-</b> -
H → WW <sup>(1)</sup> → loto vi = 7 for [Lot = 4.5 fo <sup>2</sup> vi = 6 for [Lot = 29.7 fo <sup>2</sup> ]	+
$A \rightarrow \lambda \lambda$ $A \rightarrow \lambda \lambda$	
$H \rightarrow ZZ^{(1)} \rightarrow 41$ $\dot{v}_1 - 7 \text{ TeV} \left[ (\omega_1 + 4.6 \text{ m}^2) \right]$ $\dot{v}_1 = 6 \text{ TeV} \left[ (\omega_2 - 2.2 \text{ fin}^2) \right]$	-
Combined $\mu = 1.30 \pm 0.20$ $v_{1-1} = v_{2-1} = v_{3-1} = v_{3-1}$	•
-1 (	) +1
Signal strength (μ)	

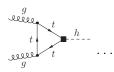
## Higgs Production and Decay in Composite Models

Higgs = Goldstone 
$$\rightarrow$$
 non-linear realization  $\Sigma_I = \left( Exp[iH_{\hat{a}}T^{\hat{a}}/f] \right)_{I5}$ 

Higgs decay constant  $f = m_{\rho}/g_{\rho} < m_{\rho}$ 

⇒ Modification of Higgs couplings due to Goldstone nature, scale as trigonometric functions of  $\boxed{v/f}$   $\Rightarrow$  indirect signs of model





$$\kappa_{p}\equivrac{g_{hpp}}{g_{hpp}^{
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### MCHM<sub>5</sub>: fermions in **5** of SO(5)

$$\kappa_f^5 = \frac{\cos(2v/f)}{\cos(v/f)} < \kappa_{W,Z}$$

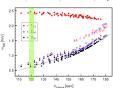
 $\kappa_{W,Z} = \cos(v/f)$ 

Giudice, Grojean, Pomarol, R. Rattazzi, hep-ph/0703164

## Effects of Resonances/Mixing? Light Top Custodians!

 Composite Higgs models (gauge-Higgs unification) feature generically light resonances associated to the (RH) top quark
 Carena, Ponton, Santiago, Wagner, hep-ph/0607106; Contino, Da Rold, Pomarol, hep-ph/0612048





• Consequence of large  $m_t$  (5D: IR localized) and enlarged fermion representations that protect  $Zb\bar{b}$  ( $P_{LR}$ )

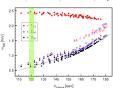
$$\zeta_R^t = \begin{bmatrix} (2,2)_R^t[-+] \\ (1,1)_R^t[+,+] \end{bmatrix}$$
$$\zeta_L^t = [+ \leftrightarrow -]$$

[+ 3 other 5s of 
$$SO(5)$$
  
 $\cong (2 \otimes 2 \oplus 1)$  of  $SU(2)_L \times SU(2)_R$ ]

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[+ 3 other **5**s of 
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New light scale  $m_{\rm cust} \ll f$  suggests that the effect of the light top custodians (incl. mixing) is dominant in these models

 Possible to describe effects due to mixing with fermion resonances in transparent way by only considering

$$\mathsf{SM} + \mathsf{light} \,\, \mathsf{custodians}$$

 Motivated vector-like fermion scenario, featuring custodial protection [neglect other effects for the moment, see later] Carmona, FG, 1301.5856

# MCHM<sub>5</sub>

$$\mathsf{IR} \; \mathsf{model:} \; \boxed{\mathsf{SM}} \; + \boxed{\mathsf{new} \; \mathsf{bi-doublets}} \subset \; \zeta_R^{t,\tau} = \left[ \begin{array}{c} (\mathbf{2},\mathbf{2})_R^{t,\tau}[-+] \\ (\mathbf{1},\mathbf{1})_R^{t,\tau}[+,+] \end{array} \right]$$

### Light top custodians:

$$Q_{1L,R}^{(0)} = egin{pmatrix} \Lambda_{1L,R}^{(0)} \ T_{1L,R}^{(0)} \end{pmatrix} \sim \mathbf{2}_{rac{7}{6}}, \quad Q_{2L,R}^{(0)} = egin{pmatrix} T_{2L,R}^{(0)} \ B_{2L,R}^{(0)} \end{pmatrix} \sim \mathbf{2}_{rac{1}{6}}$$

$$Q = 5/3, 2/3, 2/3, 1/3$$

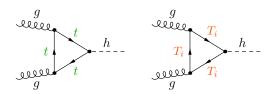
# MCHM<sub>5</sub>: Higgs Couplings of Top Sector in Mass Basis

$$g_{h5}^{T} = \frac{1}{v} \begin{pmatrix} c_{R}^{2} m_{t} & 0 & s_{R} c_{R} m_{t} \\ 0 & 0 & 0 \\ s_{R} c_{R} M_{T_{2}} & 0 & s_{R}^{2} M_{T_{2}} \end{pmatrix} \qquad \stackrel{H}{\longrightarrow} \qquad \stackrel{t}{\longrightarrow} \qquad \stackrel{T}{$$

- Mixing parameter  $1 \ge s_R^2 = 1 c_R^2 \ge 0$ : function of yukawas and vector-like masses, measures compositeness of zero mode (mixing with resonances, negligible for light quarks)
- ullet Coupling of SM mode depleted by  $c_R^2$

# Effects of Top Custodians: $gg \rightarrow h$ and $h \rightarrow \gamma \gamma$

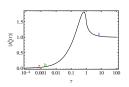
$$\sigma(gg 
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$$\kappa_g^5 pprox rac{(c_R^t)^2\,A( au_t)\ +\ (s_R^t)^2\,A( au_{T_2})}{A( au_t)} \; ; \;\;\; egin{align*} {\rm loop \ functions} \ A( au_t) pprox 1 pprox A( au_{T_2}) \ {
m neglect \ b\ (=SM) \ for \ sin} \ \end{array}$$

loop functions neglect b (=SM) for simplicity

$$\tau_i = 4m_i^2/m_h^2$$



## Effects of Top Custodians: $gg \rightarrow h$ and $h \rightarrow \gamma \gamma$

$$\sigma(gg \to h)_{\text{MCHM}_5} = |\kappa_g^5|^2 \, \sigma(gg \to h)_{\text{SM}}$$

- $\kappa_g^5 \approx (c_R^t)^2 + (s_R^t)^2 = 1$
- Effects due to quark mixing drop out after summing over SM-like top and resonances, due to  $A(\tau_t) \approx A(\tau_{T_2}) \approx 1$  see Falkowski, 0711.0828
  - $\Rightarrow$  No indirect information about details of spectrum and couplings

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- Light fermions considered elementary ⇒ no effect at first place in that sector (↔ warped models Casagrande, FG, Haisch, Neubert, Pfoh, 1005.4315;)
- Not much about fermion mixing  $(s_R)$  in past Higgs literature

$$ightarrow$$
 only  $v/f$  effects ...

Very recently: top mixing accessible in Higgs+jets Banfi, Martin, Sanz, 1308.4771

$$\Gamma(h o ff)_{
m MCHM}_5 = |\kappa_f^5|^2 \, \Gamma(h o ff)_{
m SM}$$

$$\kappa_{\gamma}^{5}=1$$
 ??

...due to same cancellations as in  $\kappa_g^5$ ?

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Indeed, top contributions cancel in the same way!

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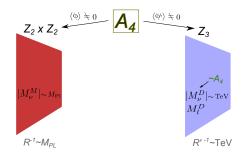
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### However, not considered before:

- $h \to \gamma \gamma$  features lepton contributions
- Negligible compositeness → no effect ?
- Not necessarily true!

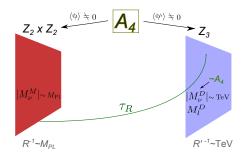
## Light Custodians: MCHM<sub>5</sub>

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- However: explaining masses and mixings with the help of (flavor protecting)  $A_4$



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- o au more composite than naively expected
- ightarrow light au custodians (setup analogous to top custodians)

del Aguila, Carmona, Santiago, 1001.5151; Csaki, Delaunay, Grojean, Grossman, 0806.0356

### Let's see what happens ...

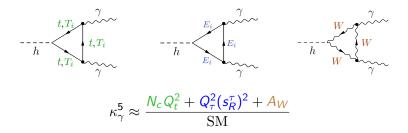
$$\kappa_{\gamma}^{5} \approx \frac{N_{c} Q_{t}^{2} + Q_{\tau}^{2}((c_{R}^{\tau})^{2} A(\tau_{\tau}) + (s_{R}^{\tau})^{2}) + A_{W}}{\text{SM}}$$

$$A_{W} \approx -6.25$$

 $A_W \approx -0.25$ dominates

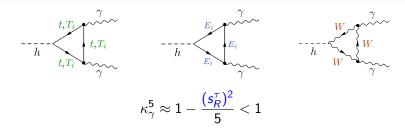
• No cancellation due to different loop fns.  $A( au_{ au}) \ll A( au_{ extsf{E}_2}) pprox 1$ 

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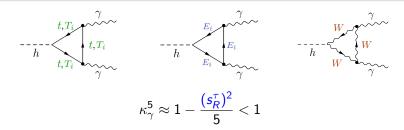
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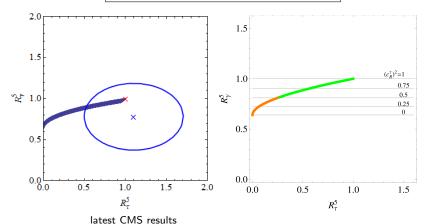


- No cancellation due to different loop fns.  $A( au_{ au}) \ll A( au_{ extsf{E}_2}) pprox 1$
- Surviving effect from fermion mixing: au custodian, compositeness of light mode au  $s_R > 0$
- More recently similar effect studied for composite light quarks see Delaunay, Grojean, Perez, 1303.5701; see also Azatov, Galloway, 1110.5646

# Higgs Phenomenology

$$\begin{array}{c} pp \rightarrow h \rightarrow \gamma\gamma \\ pp \rightarrow h \rightarrow \tau\tau \end{array}$$

$$R_f^5 \equiv \frac{[\sigma(pp \to h) {\rm Br}(h \to ff)]_{\rm MCHM_5}}{[\sigma(pp \to h) {\rm Br}(h \to ff)]_{\rm SM}}$$

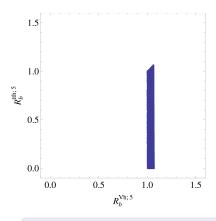


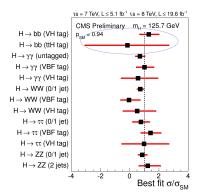
- Strong correlation allows to easily test the model
- Essentially only one parameter entering!

$$R_{\gamma}^{5} \approx \left(1 - \frac{(s_{R}^{\tau})^{2}}{5}\right)^{2}$$

$$R_{ au}^5 pprox (c_R^{ au})^4$$

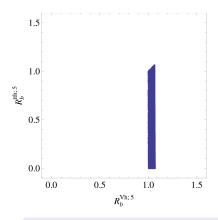
# $h \rightarrow bb$

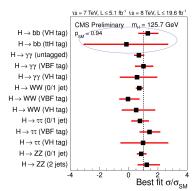




 $gg o t \bar t^* t^* \bar t o t \bar t h$ ,  $h o b \bar b$  suppressed due to  $\kappa_t^5 < 1$   $\Rightarrow$  Nice possibility to test the model in the future!

$$\sigma(tth)_{
m MCHM}_5 = (c_R^t)^4 \, \sigma(tth)_{
m SM}$$





Another option for (direct) access to parameter of model

$$(c_R^t)^2 \approx \sqrt{R_b^{tth; 5}}$$

$$\sigma(tth)_{
m MCHM_5} = (c_R^t)^4 \, \sigma(tth)_{
m SM}$$

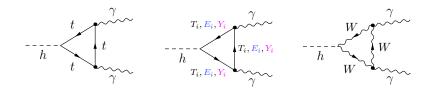
## Study Other Fermion Representations

• Put  $\tau_R$  into adjoint representation, **10** of SO(5)

$$\zeta_{R}^{\tau} = \begin{bmatrix} (\mathbf{2}, \mathbf{2})_{R}^{\tau}[-, +] & = \begin{pmatrix} N_{1R}[-, +] & E_{2R}[-, +] \\ E_{1R}[-, +] & Y_{2R}[-, +] \end{pmatrix} \\ (\mathbf{3}, \mathbf{1})_{R}^{\tau}[-, +] & = \begin{pmatrix} N_{3R}[-, +] \\ E_{3R}[-, +] \\ Y_{3R}[-, +] \end{pmatrix} \\ (\mathbf{1}, \mathbf{3})_{R}^{\tau} & = \begin{pmatrix} N_{2R}[-, +] & \tau_{R}[+, +] & Y_{1R}[-, +] \end{pmatrix} \end{bmatrix}$$

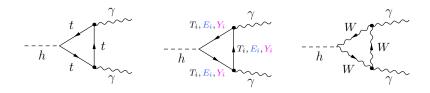
 $\mathsf{MCHM}_{5+10}$ 

## Higgs Decays: MCHM<sub>5+10</sub>



$$\kappa_{\gamma}^{5+10} \approx 1 - \frac{\nu_{E}^{5+10} + 4\nu_{Y}^{5+10}}{5} > 1$$

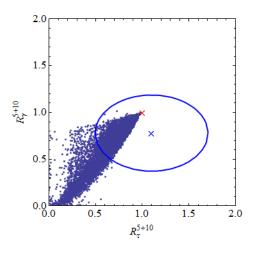
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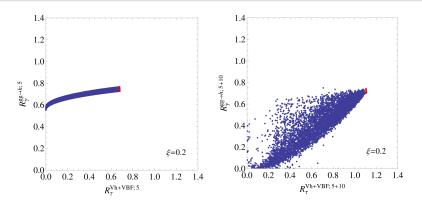
$$\kappa_{\gamma}^{5+10} \approx 1 - \frac{\nu_{E}^{5+10} + 4\nu_{Y}^{5+10}}{5} < 1$$

5D constraints on parameters  $\Rightarrow \nu_E^{5+10} > 0,\, \nu_Y^{5+10} > 0$ 

# Higgs Phenomenology: "5D" MCHM<sub>5+10</sub>

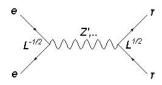


# Effects of Non-Linearity of Higgs Sector



- Trigonometric rescalings on top of fermion mixings, now also VBF and Vh production get reduced
- ullet Qualitative picture from fermion-mixing still valid in gg o h
- Usually neglected mixing effects are relevant in general

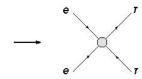
### Tau Compositeness: Estimates



Bounds on 4fermi operators  $(\bar{e} \gamma^{\mu} P_{L,R} e)(\bar{\tau} \gamma_{\mu} P_{L,R} \tau)$ :

$$C_{ee au au} < [(1-2)\,{
m TeV}]^{-2}$$

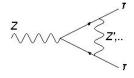
see e.g. Raidal et al., 0801.1826



In composite A4 models:

$$C_{ee\tau\tau} \sim g^2\,1/\sqrt{L}*\sqrt{L}~1/m_\rho^2$$

$$m_
ho \sim 1\,{
m TeV}$$
 ok for  $s_R^ au \sim \mathcal{O}(1)!$ 



Similar considerations  $o s_R^ au \sim {\it O}(1)$  still possible

 $\rightarrow$  exact calculation in progress

### Summary

- Lepton custodians lead to distinct phenomenology with respect to previous studies of composite models
  - $\Rightarrow$  Interesting scenario to consider
- Complementarity between direct searches for fermion partners and looking for indirect effects
- Precise measurement of Higgs couplings desirable
- Outlook: study full 5D model
- As we have seen that large signals are not to be expected from the quark sector, it could be the unexpected compositeness of the  $\tau$ -lepton that leads to first signals of compositeness in Higgs physics at the LHC

### Summary

Thank you for your attention!

### Backup: Minimal Composite Higgs Models

### Starting point for description of PGB-Higgs $(E < 4\pi f)$ :

Non-linear sigma model

$$\mathcal{L}_{\Sigma} = D_{\mu} \Sigma^{T} D^{\mu} \Sigma, \quad \Sigma_{I} = \left( \textit{Exp}[\textit{iH}_{\hat{a}} T^{\hat{a}} / f] \right)_{I5}$$

 $T^{\hat{a}}$ : (broken) generators of coset SO(5)/SO(4)

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- Expect Higgs couplings to scale as trigonometric functions of v/f
- MCHM<sub>5</sub> (MCHM<sub>10</sub>): fermions in **5** (**10**) of *SO*(5)

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 $t_R$  (residing mostly in  $(\mathbf{1},\mathbf{1})_R^u$ ) is composite  $\Rightarrow$  RH (would-be) 0-modes in  $\zeta^u$  localized moderately strong in IR  $\Rightarrow$  BCs support ultra-light KKs in  $(\mathbf{2},\mathbf{2})_R^u[-+]$ 

Contino, Da Rold, Pomarol, hep-ph/0612048 del Aguila, Carmona, Santiago, 1001.5151

[ **5** of 
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$$Q_{1L,R}^{(0)} = egin{pmatrix} \Lambda_{1L,R}^{(0)} \ T_{1L,R}^{(0)} \end{pmatrix} \sim \mathbf{2}_{rac{7}{6}}, \quad Q_{2L,R}^{(0)} = egin{pmatrix} T_{2L,R}^{(0)} \ B_{2L,R}^{(0)} \end{pmatrix} \sim \mathbf{2}_{rac{1}{6}} \ T_{R}^{3} : (rac{1}{2}, -rac{1}{2}) \end{pmatrix}$$

### Backup: MCHM<sub>5</sub>

$$\mathcal{L}_{L} = -y_{I} \bar{I}_{L}^{(0)} \varphi \tau_{R}^{(0)} - y_{I}' \Big[ \bar{L}_{1L}^{(0)} \varphi + \bar{L}_{2L}^{(0)} \tilde{\varphi} \Big] \tau_{R}^{(0)} - M_{I} \Big[ \bar{L}_{1L}^{(0)} L_{1R}^{(0)} + \bar{L}_{2L}^{(0)} L_{2R}^{(0)} \Big] + \text{h.c.}$$

$$\mathcal{L}_{Q} = -y_{q} \, \bar{q}_{L}^{(0)} \varphi t_{R}^{(0)} - y_{q}' \Big[ \bar{Q}_{1L}^{(0)} \varphi + \bar{Q}_{2L}^{(0)} \tilde{\varphi} \Big] t_{R}^{(0)} - M_{Q} \Big[ \bar{Q}_{1L}^{(0)} Q_{1R}^{(0)} + \bar{Q}_{2L}^{(0)} Q_{2R}^{(0)} \Big] + \text{h.c.}$$

$$I_L^{(0)}$$
,  $\tau_R^{(0)}$ ,  $q_L^{(0)}$ ,  $t_R^{(0)}$ : third generation SM fields,  $\varphi=1/\sqrt{2}\,(0,v+h)^T$ 

- First two generations: negligible couplings to resonances, effects of their resonances on Higgs physics negligible (different in warped XD)
- $b_R^{(0)}, \nu_R^{(0)}$  behave SM-like since there are no new resonances to which they could couple
- $P_{LR}$  symmetry:  $SU(2)_L \leftrightarrow SU(2)_R$ , protects  $Z \rightarrow b_L b_L$ ,  $Z \rightarrow \tau_R \tau_R$

# Backup: MCHM<sub>5</sub>: Spectrum

$$\mathcal{M}_{E}^{5} = \begin{pmatrix} \frac{v}{\sqrt{2}}y & 0 & 0\\ \frac{v}{\sqrt{2}}y' & M & 0\\ \frac{v}{\sqrt{2}}y' & 0 & M \end{pmatrix}$$

#### Three heavy particle with degenerate mass

### Additional heavier Q = -1 (Q = 2/3) state with

$$m_{E_2} = \frac{M}{c_R} \sqrt{1 - s_R^2 \frac{m_{\tau}^2}{M^2}}$$

### Backup: Higgs Decays

$$\Gamma(h \to ff)_{\mathrm{MCHM}_5} = |\kappa_f^5|^2 \Gamma(h \to ff)_{\mathrm{SM}}$$

$\kappa_t^5$	$(c_R^t)^2$
$\kappa_b^5$	1
$\kappa_g^5$	pprox 1
$\kappa_{ au}^{5}$	$(c_R^{\tau})^2$
$\kappa_W^5 = \kappa_Z^5$	1



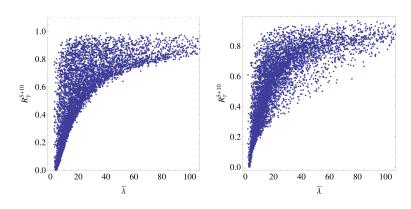
### Backup: Yukawa and Mass Lagrangian for MCHM<sub>5+10</sub>

$$\mathcal{L} = -y \, \bar{l}_{L}^{(0)} \varphi \tau_{R}^{(0)} - y' \left[ \bar{L}_{1L}^{(0)} \varphi + \bar{L}_{2L}^{(0)} \tilde{\varphi} \right] \tau_{R}^{(0)} - M \left[ \bar{L}_{1L}^{(0)} L_{1R}^{(0)} + \bar{L}_{2L}^{(0)} L_{2R}^{(0)} \right]$$

$$- \tilde{M} \left[ \bar{L}_{3L}^{(0)} L_{3R}^{(0)} + \bar{Y}_{1L}^{(0)} Y_{1R}^{0} \right] - \tilde{y} \, \bar{l}_{L}^{(0)} \sigma' \varphi L_{3R}^{(0)\prime} - \hat{y} \left[ \bar{L}_{1L}^{(0)} \sigma' \varphi - \bar{L}_{2L}^{(0)} \sigma' \tilde{\varphi} \right] L_{3R}^{(0)\prime}$$

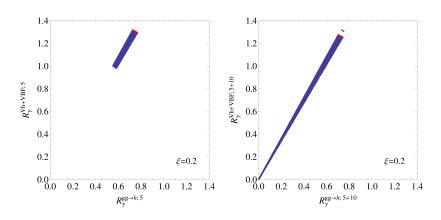
$$- \sqrt{2} \hat{y} \bar{L}_{2L}^{(0)} \varphi Y_{1R}^{(0)} + \bar{y}^* \left[ \bar{L}_{1R}^{(0)} \sigma' \varphi - \bar{L}_{2R}^{(0)} \sigma' \tilde{\varphi} \right] L_{3L}^{(0)\prime} + \sqrt{2} \bar{y}^* \bar{L}_{2R}^{(0)} \varphi Y_{1L}^{(0)} + \text{h.c.}$$

## Backup: Dependence on Parameters



$$\bar{\lambda} = \frac{2M\tilde{M}}{v^2|\bar{y}\hat{y}|}$$

# Backup: Effects of Non-Linearity of Higgs Sector



### Backup: Effects of Non-Linearity of Higgs Sector

Pseudo-Goldstone Nature of Higgs (leading order)  $\Rightarrow$ 

$$\kappa_W = \kappa_Z = \cos\left(\frac{v}{f}\right) \approx \sqrt{1-\xi}, \quad \xi = v^2/f^2$$

 $\Rightarrow$  trivial rescaling of VBF and Vh

$$\kappa_f^5 \to \kappa_f^5 \cos\left(\frac{2v}{f}\right) / \cos\left(\frac{v}{f}\right) \approx \kappa_f^5 (1 - 2\xi) / \sqrt{1 - \xi}$$

$$\kappa_g^5 \approx \cos\left(\frac{2v}{f}\right) / \cos\left(\frac{v}{f}\right) \approx (1 - 2\xi) / \sqrt{1 - \xi}$$

$$\kappa_{\tau}^{5+10} \to \kappa_{\tau}^{5+10} \cos\left(\frac{v}{f}\right) \approx \kappa_{\tau}^{5+10} \sqrt{1-\xi}$$

see Giudice, Grojean, Pomarol, R. Rattazzi, hep-ph/0703164; Azatov, Galloway, 1110.5646