

Spin-1 resonances as a signature of composite Higgs at the LHC

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with S.Pokorski, A.Weiler



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QCD Lagrangian in the limit $m_u, m_d \rightarrow 0$

$$SU(2)_L \times SU(2)_R \rightarrow SU(2)_V$$

$\sqrt{s} \ll \Lambda_{QCD}$ pions interact weakly \rightarrow effective description

$$U \rightarrow g_L U g_R^\dagger, \quad U = e^{i\pi\sigma^a/f_\pi}, \quad \mathcal{L}^{(2)} = \frac{f_\pi^2}{4} \text{Tr} \left\{ D^\mu U^\dagger D_\mu U \right\}$$

electroweak symmetry broken by new strong interactions

composite Higgs - PG boson

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composite Higgs - PG boson

- $SO(5)/SO(4) \rightarrow 4\pi \rightarrow H$

Minimal Composite Higgs Model
Agashe, Contino, Pomarol '04

- $SO(6)/SO(5) \rightarrow 5\pi \rightarrow H, a$
 $SU(4)/Sp(4, C) \rightarrow 5\pi \rightarrow H, s$

Next MCHM
Gripaios, Pomarol, Riva, Serra '09
Chacko, Batra '08

- $SO(6)/SO(4) \times SO(2) \rightarrow 8\pi \rightarrow H_1 + H_2$

Minimal Composite Two Higgs Doublets
Mrazek, Pomarol, Rattazzi, Serra, Wulzer '11

Signatures of composite Higgs

- modified Higgs couplings

$$\xi = (v/f_\pi)^2$$

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$$\mathcal{L}^{(2)} = \frac{1}{2} \partial_\mu h \partial^\mu h + \frac{v^2}{4} \left(1 + 2a \frac{h}{v} + b \frac{h^2}{v^2} + \dots \right) \text{Tr} \left\{ D^\mu U D_\mu U^\dagger \right\}$$

$$a = \sqrt{1 - \xi}, \quad b = 1 - 2\xi.$$

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→ small values of ξ preferred, $\xi \lesssim 0.22$

Signatures of composite Higgs

indirect (electroweak precision, flavor) and direct effects

- spin-1/2 resonances
- **spin-1 resonances**

Contino, Pappadopulo, Marzocca, Rattazzi

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→ analogue of ρ of QCD

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Guideline: QCD

Effective description of spin-1 resonances

global symmetry breaking $\mathcal{G} \rightarrow \mathcal{H}$

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Spin-1 resonances

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- $W, Z \leftrightarrow \rho$ mixing
gauge kinetic terms \rightarrow interactions of spin-1 resonances with W, Z eigenstates
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- *fermion* \leftrightarrow *fermion partner* mixing
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- modify the symmetry breaking pattern

$$\mathcal{G} \times \mathcal{H}_{local} \rightarrow \mathcal{H}$$

- SM electroweak $SU(2)_L \times U(1)_Y$ group sits in \mathcal{G}
- gauge bosons of \mathcal{H}_{local} → 'vector' resonances

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$$S \rightarrow g S h^\dagger, \quad g \in \mathcal{G}, \quad h \in \mathcal{H}_{local}, \quad \langle S \rangle = \mathbf{1}.$$

$$\mathcal{L} \ni v_1^2 \text{Tr} \left\{ D_\mu S D^\mu S^\dagger \right\}$$

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at leading order in g/g_ρ

- heavy spin-1 eigenstates \leftrightarrow 'hidden gauge' ρ^μ fields
- light eigenstates \leftrightarrow SM A, W, Z fields
- mixing $\sim g, g'/g_\rho \rightarrow$ interactions!

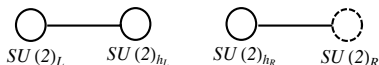
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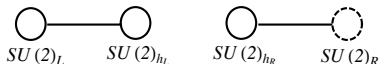
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- 3 free parameters: $\xi, g_\rho, g_{\rho\pi\pi}$

$$g_{\rho\pi\pi} \epsilon^{abc} \pi^a \partial_\mu \pi^b \rho_\mu^c - g_\rho \epsilon^{abc} \partial_\mu \rho_\nu^a \rho_\mu^b \rho_\nu^c$$

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$$W_\mu^\pm \approx \tilde{W}_\mu^\pm - \frac{\sqrt{2}}{2} \sqrt{2-\xi} \frac{g}{g_\rho} \tilde{\rho}_{L\mu}^\pm$$

$$Z_\mu \approx \tilde{Z}_\mu - \frac{\sqrt{2-\xi}}{\sqrt{2}} \frac{g^2 - g'^2}{g_\rho \sqrt{g^2 + g'^2}} \tilde{\rho}_{L\mu}^0 - \frac{2\sqrt{2-2\xi}}{(2-\xi)^{3/2}} \frac{g'^2}{g_\rho \sqrt{g^2 + g'^2}} \tilde{\rho}_{R\mu}^0$$

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assumption: couplings of $\tilde{\rho}$ eigenstates with SM fermions arise only via their admixture in SM W_μ^\pm, Z_μ

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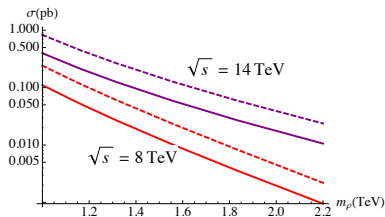
- coupling of ρ to two fermions enhanced for small ξ
- coupling of ρ to two SM gauge bosons suppressed

$$g_{\rho\pi\pi} = \xi \frac{m_\rho^2}{\sqrt{2}g_\rho v^2} = \frac{m_\rho^2}{\sqrt{2}g_\rho f_\pi^2}$$

Production and decays

- production dominated by Drell-Yan $q\bar{q} \rightarrow \rho$

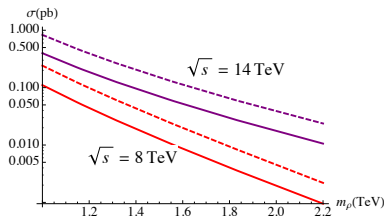
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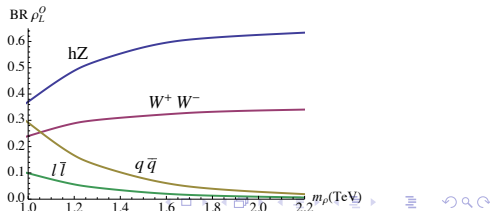
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- decays mainly to hZ and WW , but ll non-negligible

for a specific value of $\xi = 0.2$ and $g_\rho = 4$

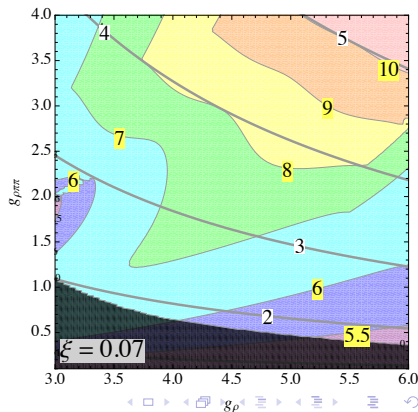
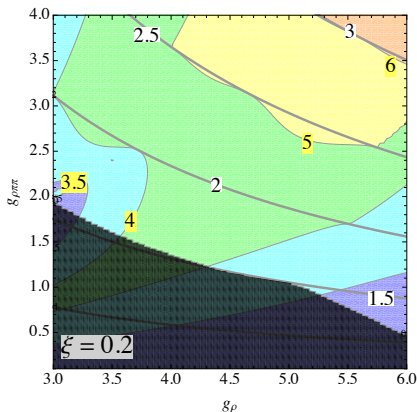
$$\Gamma(\rho \rightarrow WW) \sim m_\rho g_\rho^2 \pi$$



Direct searches

most sensitive: CMS search for dilepton resonances

$$m_\rho^2 \approx \frac{\sqrt{2}g_\rho g_{\rho\pi\pi} v^2}{\xi}$$



Conclusions

- signatures of composite Higgs - modified Higgs couplings, effects of resonances
- general effective framework for spin-1 resonances → phenomenology
- at small ξ - the spin-1 resonance coupling to two SM gauge bosons is suppressed, the coupling to two fermions is enhanced
- resonances mainly Drell-Yan produced
- exclusion limits from searches for dilepton resonances, diboson resonances, dijet mass spectra, ...
- most sensitive now - searches for dilepton resonances, $m_\rho \lesssim 2\text{TeV}$ excluded

Perturbative unitarity constraints

without spin-1 resonances $\mathcal{M}_{WW \rightarrow WW}^0(s) \sim \frac{1}{16\pi} \frac{\xi s}{v^2} = \frac{1}{16\pi} \frac{s}{f_\pi^2}$

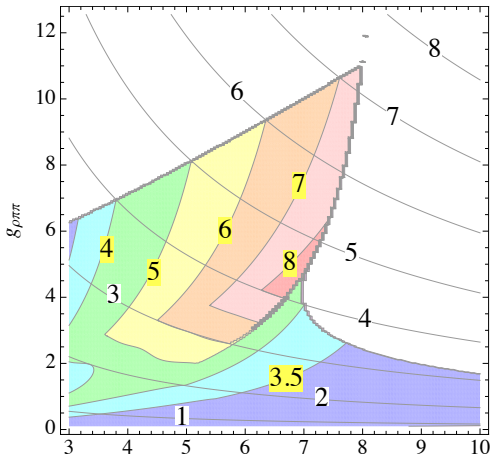
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→ perturbative unitarity violation at $\Lambda \sim 1.3 \text{ TeV}/\sqrt{\xi}$

$$\xi = 0.2$$

add ρ_L and ρ_R
resonances,
inelastic channels
included

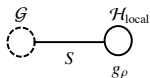
$$m_\rho^2 \approx \frac{2g_\rho g_{\rho\pi\pi} v^2}{\xi}$$



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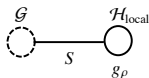
- 'vector' resonances $\mathcal{G} \times \mathcal{H}_{local} \rightarrow \mathcal{H}$



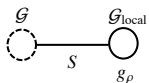
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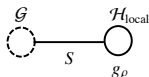
- 'vector' and 'axial' resonances $\mathcal{G} \times \mathcal{G}_{local} \rightarrow \mathcal{H}$



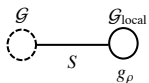
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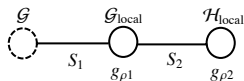
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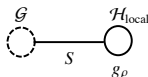
- more resonances $\mathcal{G} \times \mathcal{G}_{local} \times \mathcal{H}_{local} \rightarrow \mathcal{H}$



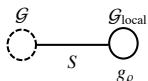
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- 'vector' resonances most relevant for phenomenology