Towards Determining the UV Behavior of Maximal Supergravity

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Based on work with: Zvi Bern, John Joseph Carrasco, Lance Dixon, Radu Roiban

Text-book: perturbative gravity is complicated !

de Donder gauge:
$$\mathcal{L}=rac{2}{\kappa^2}\sqrt{g}R, \quad g_{\mu
u}=\eta_{\mu
u}+\kappa h_{\mu
u}$$

$$\sum_{\mu_1}^{\nu_1} \sum_{\mu_2}^{\nu_2} = \frac{1}{2} \left[\eta_{\mu_1\nu_1} \eta_{\mu_2\nu_2} + \eta_{\mu_1\nu_2} \eta_{\nu_1\mu_2} - \frac{2}{D-2} \eta_{\mu_1\mu_2} \eta_{\nu_1\nu_2} \right] \frac{i}{p^2 + i\epsilon}$$

$$\begin{array}{l} k_{2} \\ \mu_{2} \\ \mu_{2} \\ \mu_{2} \\ \mu_{3} \\ \mu_{4} \\ \mu_{4} \\ \mu_{1} \\ k_{1} \\ \mu_{1} \end{array} = \operatorname{sym} \begin{bmatrix} -\frac{1}{2} P_{3}(k_{1} \cdot k_{2}\eta_{\mu_{1}\nu_{1}}\eta_{\mu_{2}\nu_{2}}\eta_{\mu_{3}\nu_{3}}) - \frac{1}{2} P_{6}(k_{1\mu_{1}}k_{1\nu_{2}}\eta_{\mu_{1}\nu_{1}}\eta_{\mu_{3}\nu_{3}}) + \frac{1}{2} P_{3}(k_{1} \cdot k_{2}\eta_{\mu_{1}\mu_{2}}\eta_{\nu_{1}\nu_{2}}\eta_{\mu_{3}\nu_{3}}) \\ + P_{6}(k_{1} \cdot k_{2}\eta_{\mu_{1}\nu_{1}}\eta_{\mu_{2}\mu_{3}}\eta_{\nu_{2}\nu_{3}}) + 2P_{3}(k_{1\mu_{2}}k_{1\nu_{3}}\eta_{\mu_{1}\nu_{1}}\eta_{\nu_{2}\mu_{3}}) - P_{3}(k_{1\nu_{2}}k_{2\mu_{1}}\eta_{\nu_{1}\mu_{1}}\eta_{\mu_{3}\nu_{3}}) \\ + P_{3}(k_{1\mu_{3}}k_{2\nu_{3}}\eta_{\mu_{1}\mu_{2}}\eta_{\nu_{1}\nu_{2}}) + P_{6}(k_{1\mu_{3}}k_{1\nu_{3}}\eta_{\mu_{1}\mu_{2}}\eta_{\nu_{1}\nu_{2}}) + 2P_{6}(k_{1\mu_{2}}k_{2\nu_{3}}\eta_{\nu_{2}\mu_{1}}\eta_{\nu_{1}\mu_{3}}) \\ + 2P_{3}(k_{1\mu_{2}}k_{2\mu_{1}}\eta_{\nu_{2}\mu_{3}}\eta_{\nu_{3}\nu_{1}}) - 2P_{3}(k_{1} \cdot k_{2}\eta_{\nu_{1}\mu_{2}}\eta_{\nu_{2}\mu_{3}}\eta_{\nu_{3}\mu_{1}})] \\ After symmetrization \\ \sim 100 \text{ terms }! \end{array}$$

higher order vertices...



Sh .

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On-shell simplifications

Graviton plane wave:

$$\varepsilon^{\mu}(p)\varepsilon^{\nu}(p)\,e^{ip\cdot x}$$

Yang-Mills polarization

On-shell 3-graviton vertex:

$$\sum_{\mu_{2}}^{k_{2}} \sum_{\nu_{3}}^{\nu_{2}} \sum_{\nu_{3}}^{\mu_{3}} k_{3} = i\kappa \Big(\eta_{\mu_{1}\mu_{2}}(k_{1}-k_{2})_{\mu_{3}} + \text{cyclic} \Big) \Big(\eta_{\nu_{1}\nu_{2}}(k_{1}-k_{2})_{\nu_{3}} + \text{cyclic} \Big)$$

$$\sum_{\mu_{1}}^{\nu_{1}} \sum_{\mu_{1}}^{\mu_{3}} \sum_{\nu_{3}}^{\mu_{3}} k_{3} = i\kappa \Big(\eta_{\mu_{1}\mu_{2}}(k_{1}-k_{2})_{\mu_{3}} + \text{cyclic} \Big) \Big(\eta_{\nu_{1}\nu_{2}}(k_{1}-k_{2})_{\nu_{3}} + \text{cyclic} \Big)$$

$$\sum_{\mu_{1}}^{\nu_{1}} \sum_{\mu_{1}}^{\mu_{1}} \sum_{\nu_{3}}^{\mu_{3}} k_{3} = i\kappa \Big(\eta_{\mu_{1}\mu_{2}}(k_{1}-k_{2})_{\mu_{3}} + \text{cyclic} \Big) \Big(\eta_{\nu_{1}\nu_{2}}(k_{1}-k_{2})_{\nu_{3}} + \text{cyclic} \Big) \Big) \Big(\eta_{\nu_{1}\nu_{2}}(k_{1}-k_{2})_{\nu_{3}} + \text{cyclic} \Big) \Big(\eta_{\nu_{1}\nu_{2}}(k_{1}-k_{2})_{\nu_{3}} + \text{cyclic} \Big) \Big) \Big(\eta_{\nu_{1}\nu_{3}}(k_{1}-k_{2})_{\nu_{3}} + \frac{\eta_{1}\nu_{2}}(k_{1}-k_{2})_{\nu_{3}} + \frac{\eta_{1}\nu_{2}}(k_{1}-k_{2})_{\nu_{3}} + \frac{\eta_{1}\nu_{2}}(k_{1}-k_{2})_{\nu_{3}} + \frac{\eta_{1}\nu_{2}}(k_{1}-k_{2})_{\nu_{3}} + \frac{\eta_{1}\nu_{2}}(k_{1}-k_{2})_{\nu_{3}} + \frac{\eta_{1}\nu_{2}}(k_{1}-k_{2})_{\nu_{3}} + \frac{\eta_{1}\nu_{2}}(k_{1}-k_{$$

Gravity scattering amplitude:

$$\mathcal{M}_{\text{tree}}^{\text{GR}}(1,2,3,4) = \frac{st}{u} A_{\text{tree}}^{\text{Yang-Mills amplitude}} (1,2,3,4) \otimes A_{\text{tree}}^{\text{YM}}(1,2,3,4)$$

Gravity processes = squares of gauge theory ones - entire S-matrix Bern, Carrasco, HJ [BCJ]

Gravity should be simple off shell



UV problem = basic power counting

Naively expect gravity to behave worse than Yang-Mills

Gravity: non-renormalizable dimensionful coupling



$$d^{4L}p \frac{\dots (\kappa p^{\mu}p^{\nu})\dots}{p_1^2 p_2^2 p_3^2 \dots p_n^2}$$

 $\sim \int d^{4L} p \frac{\dots (gp^{\mu}) \dots}{p_1^2 p_2^2 p_2^2 \dots p_2^2}$

momenta

Yang-Mills: renormalizable dimensionless coupling

For finite gravity \rightarrow vast cancellations needed $\sim (p^{\mu})^{2L} \rightarrow (k^{\mu})^{2L}$ seems implausible, but exists for *N*=8 SG in all known ampl's.

Outline

UV status of N=8 SUGRA

- Duality between color and kinematics
 - Double-copy structure of gravity
 - Ability to calculate
- Amplitude UV behavior from duality
- Current 5-loop SUGRA progress
- Conclusion

SUGRA status on one page

Facts:

- No D=4 divergence of pure SG has been found to date.
- Susy forbids 1,2 loop div. R^3 , R^3

Ferrara, Zumino, Deser, Kay, Stelle, Howe, Lindström, Green, Schwarz, Brink, Marcus, Sagnotti

- Pure gravity 1-loop finite, 2-loop divergent Goroff & Sagnotti
- With matter: 1-loop divergent 't Hooft & Veltman
- Naively susy allows 3-loop div. R⁴
- N=8 SG and N=4 SG 3-loop finite! Bern, Carrasco, Dixon, HJ, Kosower, Roiban, Davies, Dennen, Huang
- \square $\mathcal{N}=8$ SG: no divergence before 7 loops
- D>4 divergences obey: $D_c = 4 + \frac{6}{L}$ (L>1) Marcus, Sagnotti, Bern, Dixon, Dunbar, Perelstein, Rozowsky, Carrasco, HJ, Kosower, Roiban
- 7-loop div. in D=4 implies a 5-loop div. in D=24/5 – calculation in progress!



Why is it interesting?

- If $\mathcal{N}=8$ SG is perturbatively finite, why is it interesting?
- It better be finite for a good reason!
 - Hidden new symmetry, for example
 - Understanding the mechanism might open a host of possibilities
- Any indication of hidden structures yet?
 - Gravity is a double copy of gauge theories
 - Color-Kinematics: kinematics = Lie algebra Bern, Carrasco, HJ
 - Constraints from E-M duality ? Kallosh,....
 - Hidden superconformal N=4 SUGRA ? Ferrara, Kallosh, Van Proeyen
 - Extended N=4 superspace ? Bossard, Howe, Stelle



Known UV divergences in D>4

Plot of critical dimensions of \mathcal{N} = 8 SUGRA and \mathcal{N} = 4 SYM



Known UV divergences in D>4

Plot of critical dimensions of \mathcal{N} = 8 SUGRA and \mathcal{N} = 4 SYM



L = 7 lowest loop order for possible D = 4 divergence

Beisert, Elvang, Freedman, Kiermaier, Morales, Stieberger; Björnsson, Green, Bossard, Howe, Stelle, Vanhove Kallosh, Ramond, Lindström, Berkovits, Grisaru,

Siegel, Russo, Cederwall, Karlsson, and more....

Known UV divergences in D>4

Plot of critical dimensions of \mathcal{N} = 8 SUGRA and \mathcal{N} = 4 SYM



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Historical record – where is the \mathcal{N} = 8 div. ?

3 loops	Conventional superspace power counting	Green, Schwarz, Brink (1982) Howe and Stelle (1989) Marcus and Sagnotti (1985)
5 loops	Partial analysis of unitarity cuts; If \mathcal{N} = 6 harmonic superspace exists; algebraic renormalisation	Bern, Dixon, Dunbar, Perelstein, Rozowsky (1998) Howe and Stelle (2003,2009)
6 loops	If \mathcal{N} = 7 harmonic superspace exists	Howe and Stelle (2003)
7 loops	If \mathcal{N} = 8 harmonic superspace exists; string theory U-duality analysis; lightcone gauge locality arguments; $E_{7(7)}$ analysis, unique 1/8 BPS candidate	Grisaru and Siegel (1982); Green, Russo, Vanhove; Kallosh; Beisert, Elvang, Freedman, Kiermaier, Morales, Stieberger; Bossard, Howe, Stelle, Vanhove
8 loops	Explicit identification of potential susy invariant counterterm with full non-linear susy	Howe and Lindström; Kallosh (1981)
9 loops	Assume Berkovits' superstring non-renormalization theorems can be carried over to \mathcal{N} = 8 supergravity	Green, Russo, Vanhove (2006)
Finite	Identified cancellations in multiloop amplitudes; lightcone gauge locality and $E_{7(7)}$, inherited from hidden N=4 SC gravity	Bern, Dixon, Roiban (2006), Kallosh (2009–12), Ferrara, Kallosh, Van Proeyen (2012)

note: above arguments/proofs/speculation are only lower bounds

 \rightarrow only an explicit calculation can prove the existence of a divergence!

Color-Kinematics Duality

Color-Kinematics Duality

Yang-Mills theories are controlled by a kinematic Lie algebra

• Amplitude represented by cubic graphs:

$$\mathcal{A}_{m}^{(L)} = \sum_{i \in \Gamma_{3}} \int \frac{d^{LD}\ell}{(2\pi)^{LD}} \frac{1}{S_{i}} \frac{n_{i}c_{i}}{p_{i_{1}}^{2}p_{i_{2}}^{2}p_{i_{3}}^{2}\cdots p_{i_{l}}^{2}} \leftarrow \text{ propagators}$$
Color & kinematic numerators satisfy same relations:

$$\int \frac{d^{LD}\ell}{(2\pi)^{LD}} \frac{1}{S_{i}} \frac{n_{i}c_{i}}{p_{i_{1}}^{2}p_{i_{2}}^{2}p_{i_{3}}^{2}\cdots p_{i_{l}}^{2}} \leftarrow \text{ propagators}$$
Jacobi identity
$$\int \frac{adc}{f^{ceb}} = \int \frac{eac}{f^{cbd}} - \int \frac{abc}{f^{cde}} \int \frac{d^{LD}\ell}{f^{bac}} = \int \frac{eac}{f^{cbd}} - \int \frac{f^{abc}}{f^{abc}}$$

Duality: color ↔ kinematics

Bern, Carrasco, HJ

Some details of color-kinematics duality

Bern, Carrasco, HJ

can be checked for 4pt on-shell ampl. using Feynman rules

Example with two quarks:



- **1.** $(A^{\mu})^4$ contact interactions absorbed into cubic graphs
 - by hand 1=s/s
 - or by auxiliary field $B \sim (A^{\mu})^2$
- 2. Beyond 4-pts duality not automatic \rightarrow Lagrangian reorganization
- 3. Known to work at tree level: all-*n* example Kiermaier; Bjerrum-Bohr et al.
- 4. Enforces (BCJ) relations on partial amplitudes \rightarrow (*n*-3)! basis
- 5. Same/similar relations control string theory S-matrix

Bjerrum-Bohr, Damgaard, Vanhove; Stieberger

Gravity is a double copy

• Gravity amplitudes obtained by replacing color with kinematics

$$\mathcal{A}_{m}^{(L)} = \sum_{i \in \Gamma_{3}} \int \frac{d^{LD}\ell}{(2\pi)^{LD}} \frac{1}{S_{i}} \frac{n_{i}c_{i}}{p_{i_{1}}^{2}p_{i_{2}}^{2}p_{i_{3}}^{2}\cdots p_{i_{l}}^{2}}$$
BCJ
$$\mathcal{M}_{m}^{(L)} = \sum_{i \in \Gamma_{3}} \int \frac{d^{LD}\ell}{(2\pi)^{LD}} \frac{1}{S_{i}} \frac{n_{i}\tilde{n}_{i}}{p_{i_{1}}^{2}p_{i_{2}}^{2}p_{i_{3}}^{2}\cdots p_{i_{l}}^{2}}$$

• The two numerators can belong to different theories:

$$\begin{array}{cccc} n_i & \tilde{n}_i \\ (\mathcal{N}=4) \times (\mathcal{N}=4) & \rightarrow & \mathcal{N}=8 \text{ sugra} \\ (\mathcal{N}=4) \times (\mathcal{N}=2) & \rightarrow & \mathcal{N}=6 \text{ sugra} \\ (\mathcal{N}=4) \times (\mathcal{N}=0) & \rightarrow & \mathcal{N}=4 \text{ sugra} \\ (\mathcal{N}=4) \times (\text{matter}) & \rightarrow & \mathcal{N}=4 \text{ matter} \\ (\mathcal{N}=0) \times (\mathcal{N}=0) & \rightarrow & \text{Einstein gravity + axion+ dilaton} \end{array}$$

UV calculations using C-K duality

C-K duality unifies 1,2-loop UV behavior

1,2 loops N=4,5,6,8 SG are particularly easy to understand



But N=4 matter diverges:
Fischler; Fradkin, TseytlinImage: Comparison of the set of the set

Bern, Davies, Dennen, Huang 1209.2472 [hep-th]

3-loop \mathcal{N} =8 SG & \mathcal{N} =4 SYM

Color-kinematics dual form:

Bern, Carrasco, HJ

$$N^{(e)} = s(\tau_{45} + \tau_{15}) + \frac{1}{3}(t-s)(s+\tau_{15} - \tau_{25})$$

-







 $au_{ij} = 2k_i \cdot l_j$

UV divergent in *D*=6: Bern, Carrasco, Dixon, HJ, Roiban

$$\mathcal{A}^{(3)}\Big|_{\text{pole}} = 2g^8 st A^{\text{tree}} (N_c^3 V^{(A)} + 12N_c (V^{(A)} + 3V^{(B)})) \times (u \text{Tr}[T^{a_1} T^{a_2} T^{a_3} T^{a_4}] + \text{ perms})$$

$$\mathcal{M}^{(3)}\Big|_{\text{pole}} = 10 \Big(\frac{\kappa}{2}\Big)^8 (stu)^2 M^{\text{tree}} (V^{(A)} + 3V^{(B)})$$





4-loops \mathcal{N} =4 SYM and \mathcal{N} =8 SG

Bern, Carrasco, Dixon, HJ, Roiban 1201.5366

- 85 diagrams
- Power counting manifest
- • \mathcal{N} =4 & \mathcal{N} =8 diverge in D=11/2

$$N_{6}^{\text{SYM}} = \frac{1}{2} s_{12}^{2} (\tau_{45} - \tau_{35} - s_{12})$$

$$N_{6}^{\text{SG}} = \left[\frac{1}{2} s_{12}^{2} (\tau_{45} - \tau_{35} - s_{12})\right]^{2} \xrightarrow{2}_{1} - \underbrace{4}_{6} \underbrace{5}_{5} \underbrace{5}_{7} \underbrace{4}_{4}$$

$$\tau_{ij} = 2k_{i} \cdot l_{j}$$

- 🙊

$$\mathcal{A}_{4}^{(4)}\Big|_{\text{pole}} = -6g^{10}stA^{\text{tree}}N_{c}^{2}\left(N_{c}^{2}V_{1} + 12(V_{1} + 2V_{2} + V_{8})\right) \times (u\text{Tr}_{1234} + \text{perms})$$

$$\mathcal{M}_4^{(4)}\Big|_{\text{pole}} = -\frac{23}{8} \left(\frac{\kappa}{2}\right)^{10} stu(s^2 + t^2 + u^2)^2 M^{\text{tree}}(V_1 + 2V_2 + V_8)$$

up to overall factor, divergence same as for $\mathcal{N}=4$ SYM $1/N_c^2$ part



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5-loop status

\mathcal{N} =4 SYM 5-loop Amplitude

 $\mathcal{N}\text{=}4$ SYM important stepping stone to $\mathcal{N}\text{=}8$ SG

1207.6666 [hep-th] Bern, Carrasco, HJ, Roiban

• 416 integral topologies: (335)(370)(404)(410) Used maximal cut method Bern, Carrasco, HJ, Kosower MC NMC N^2MC N³MC • UV divergence in D=26/5: (b) (c) $\mathcal{A}_{4}^{(5)}\Big|_{\text{div}} = -\frac{144}{5}g^{12}stA_{4}^{\text{tree}}N_{c}^{3}\left(N_{c}^{2}V^{(a)} + 12(V^{(a)} + 2V^{(b)} + V^{(c)})\right)$ $\times \mathrm{Tr}[T^{a_1}T^{a_2}T^{a_3}T^{a_4}]$ H. Johansson SUSY 2013

UV calc. makes it to Hollywood





The Parking Spot Escalation

with some help from Sheldon Cooper...

\mathcal{N} =8 SG 5-loop Status

Construction using only unitarity difficult

- Works for 5-loop *N*=4 SYM
- 5-loop SG too difficult this way (ansatz: billions of terms)



 $n^{\rm SYM} \sim 8000 \text{ terms}$ $n^{\rm SG} \sim (8000)^2/2$ $\sim 30\,000\,000 \text{ terms}$

Only way: use color-kinematics duality

In principle: need only reorganize 5-loop $\mathcal{N}=4$ SYM Bern, Carrasco, HJ, Roiban

- 416 + 336 = 752 integral topologies
- Minimal ansatz: **1112** free parameters
- 2500 functional Jacobi eqns $\,\sim\,$ 20 000 000 linear eqns
- Solution exists: 29 free parameters
- Unfortunately, not all unitarity cuts work, there's some glitch
- Working with enlarged Ansätze ... stay tuned for results!



Summary

- Explicit calculations in $\mathcal{N} = 8$ SUGRA up to four loops show that the power counting exactly follows that of $\mathcal{N} = 4$ SYM a finite theory
- 5 loop calculation in D=24/5 probes the potential 7-loop D=4 counterterm will provide critical input to the $\mathcal{N}=8$ question !
- Color-Kinematics duality allows for gravity calculations for multiloop multipoint amplitudes – greatly facilitating UV analysis in gravity.
- Numbers in UV divergences of $\mathcal{N}=8$ SUGRA and $1/N_c^2$ $\mathcal{N}=4$ SYM coincide, suggesting a deeper connection between the theories
- Stay tuned for the 5-loop SUGRA result...



