

Soft R-Parity Violation

Yuhsin Tsai



in collaboration with **Gordan Krnjaic**

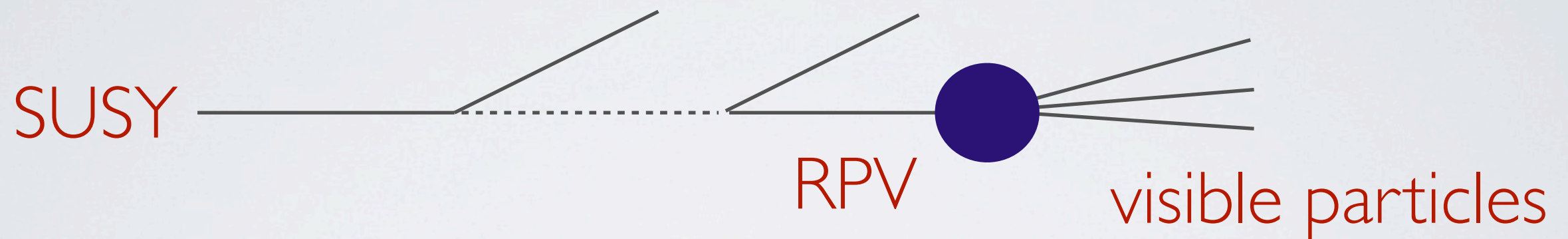


arXiv : 1304.7004

Trieste , SUSY 2013

R-parity violation (RPV)

The LSP decay fails the MET searches



However, there are stringent flavor constraints!

proton decay , neutron/anti-neutron oscillations ,
di-nucleon decay , FCNC

Soft RPV model

Generate RPV couplings through
soft SUSY breaking terms

Motivations

embed RPV into a SUSY breaking setup,
make the model more economic

A SoftRPV model needs...

- a symmetry G that forbids the supersymmetric RPV couplings
- G -breaking by SUSY breaking soft terms in a hidden sector
- mediation of the G -breaking effect to visible sector

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$$\cancel{W_{RPV}} = \frac{\lambda_{ijk}}{2} L_i L_j \bar{E}_k + \lambda'_{ijk} Q_i L_j \bar{D}_k + \frac{\lambda''_{ijk}}{2} \bar{U}_i \bar{D}_j \bar{D}_k + \mu_{L_i} L_i H_u$$

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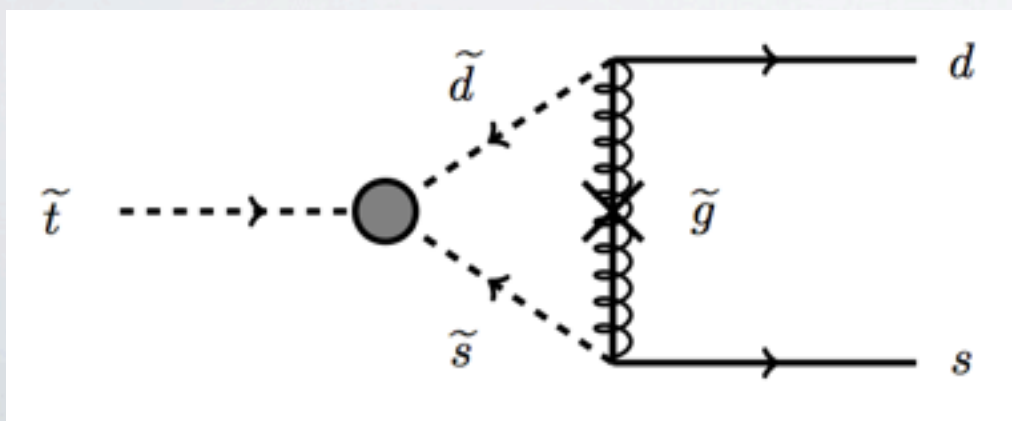
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$$\mathcal{A}_{u_i d_j d_k} \tilde{U}_i \tilde{D}_j \tilde{D}_k$$

Naively, $\mathcal{A}_{u_i d_j d_k} \simeq \mathcal{A}_{soft}$

but $\lambda''_{udd} \simeq \frac{g_s^2}{16\pi^2} \frac{\mathcal{A}_{soft}}{m_{\tilde{g}}} \simeq 10^{-2}$

which is forbidden by L- or B-violating constraints



A Gauge-Mediation Example

In this example

embed RPV in a GM setup

use R-symmetry to forbid the supersymmetric RPV terms

gravitino serves as a good dark matter candidate

no extra mass scales or small couplings need to be put in by hand

$$\lambda''_{udd} \sim \frac{g_s^2}{16 \pi^2} \frac{m_{soft}}{\sqrt{F}}$$

To forbid the supersymmetric RPV

use the R-symmetry to forbid the RPV couplings

assign

$$R[Q, \bar{U}, \bar{D}] = 1, \quad R[L] = 4/3, \quad R[\bar{E}] = 2/3, \quad R[H_u, H_d] = 0$$

$$W_{RPV} = \frac{\lambda_{ijk}}{2} \cancel{L_i L_j \bar{E}_k} + \lambda'_{ijk} \cancel{Q_i L_j \bar{D}_k} + \frac{\lambda''_{ijk}}{2} \cancel{\bar{U}_i \bar{D}_j \bar{D}_k} + \mu_I \cancel{L_i H_u}$$

the fractional R-charges are used to forbid the generation of lepton-related operators

Breaking & mediation

Gauge mediation model

~~SUSY~~

$X \Sigma \bar{\Sigma}$

soft R-breaking sector

R-symmetry

$\bar{U} \bar{D}$

$Q L \bar{E} H_u H_d$

visible sector

No RPV couplings

Breaking & mediation

Gauge mediation model

~~SUSY~~

R-symmetry

$X \quad \Sigma \quad \bar{\Sigma}$

soft R-breaking sector

$\bar{U} \quad \bar{D}$
 $Q \quad L \quad \bar{E} \quad H_u \quad H_d$

visible sector

No RPV couplings

Breaking & mediation

Gauge mediation model

~~SUSY~~

$X \quad \Sigma \quad \bar{\Sigma}$

soft R-breaking sector

~~R-symmetry~~

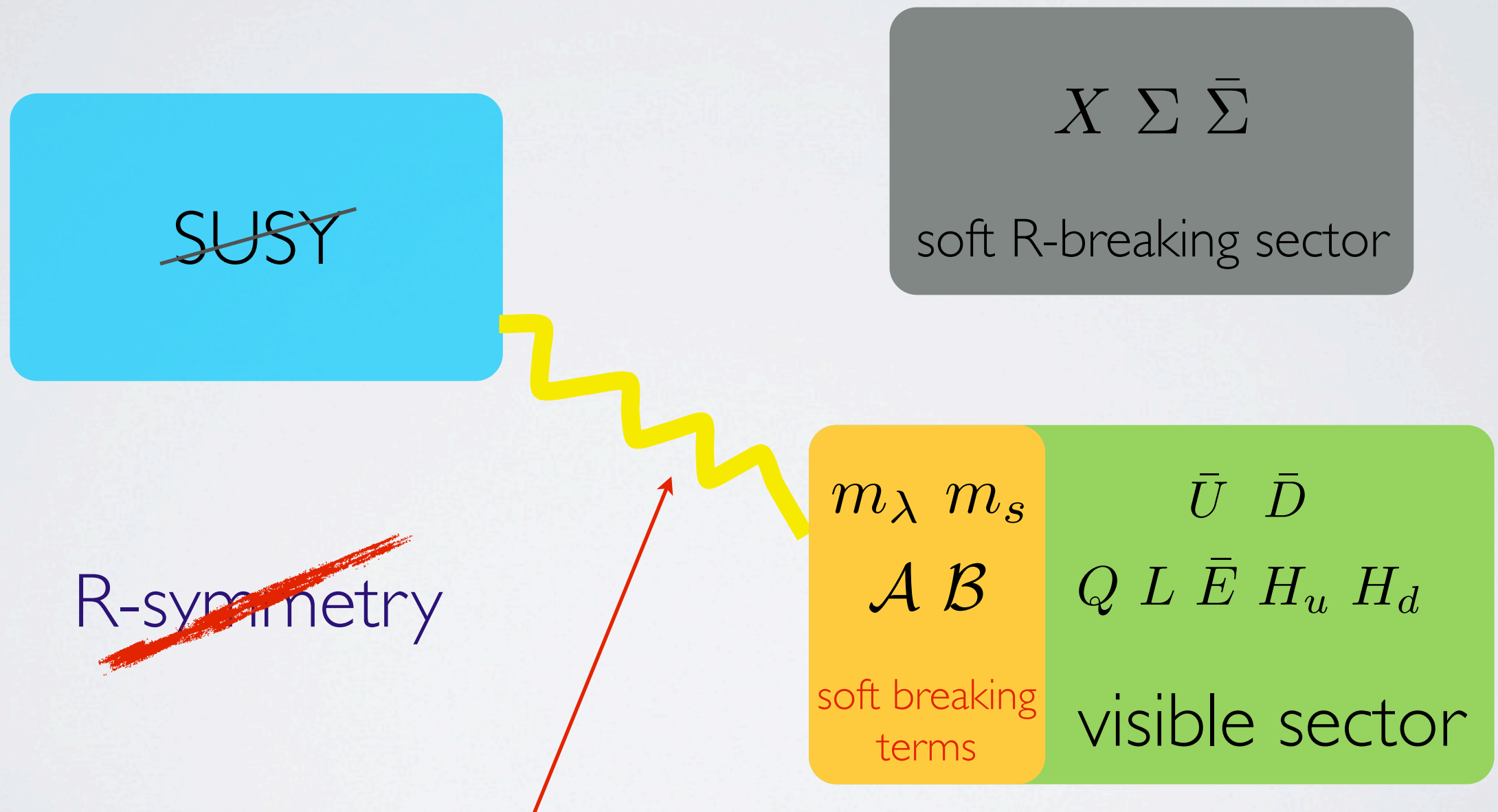
$\bar{U} \quad \bar{D}$
 $Q \quad L \quad \bar{E} \quad H_u \quad H_d$

visible sector

No RPV couplings

Breaking & mediation

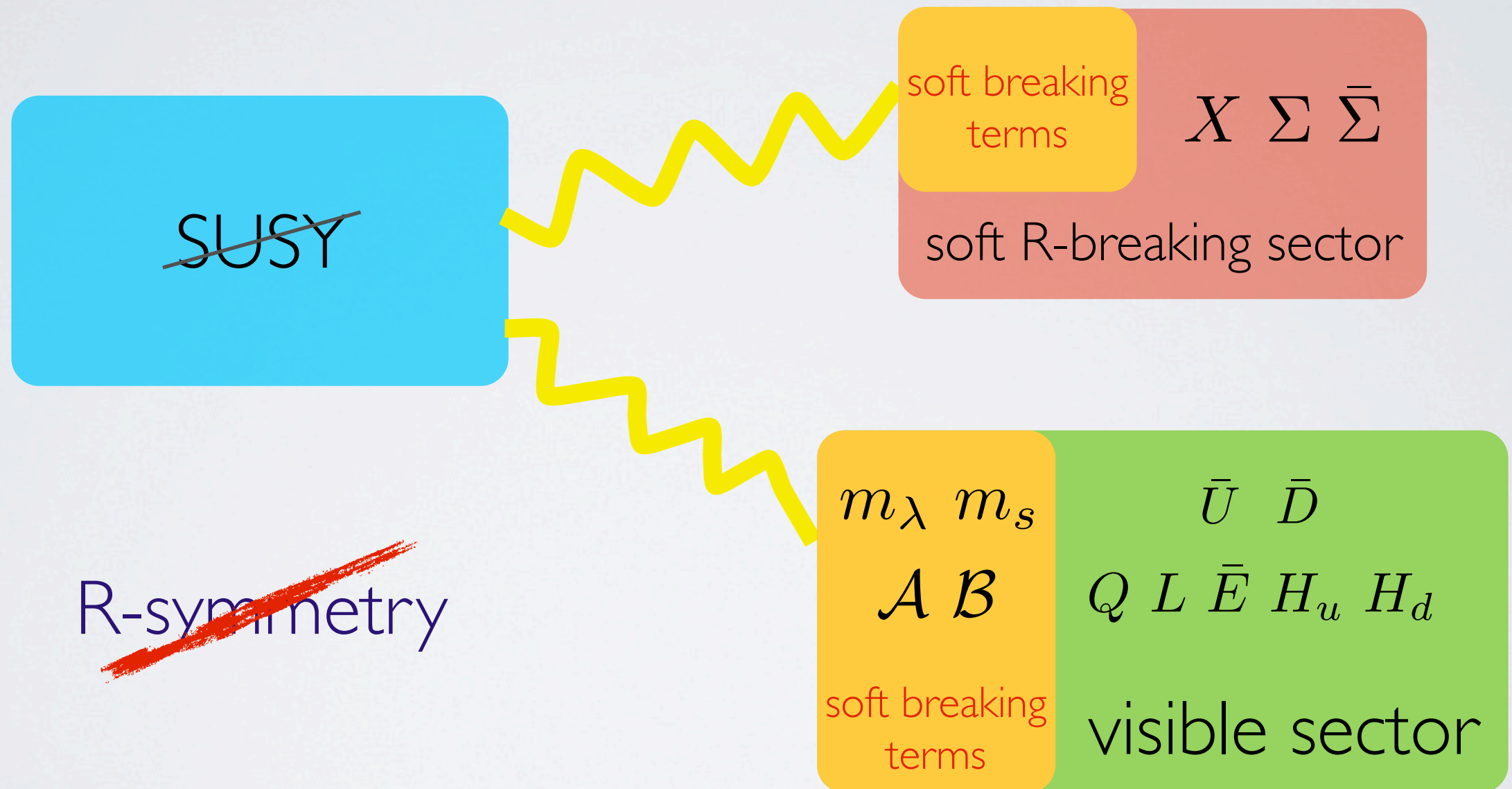
Gauge mediation model



No B- or L-violation, still no RPV couplings

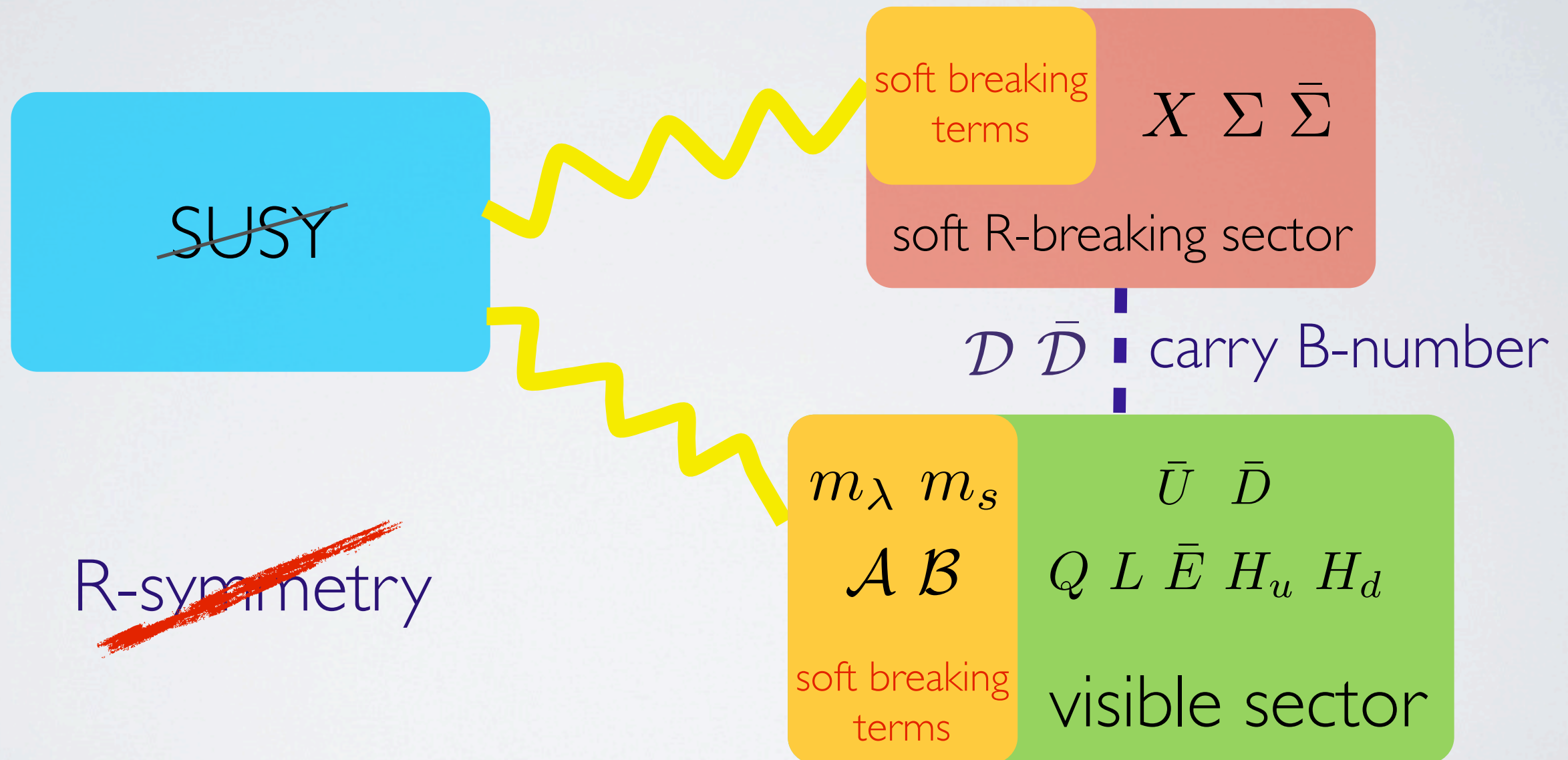
Breaking & mediation

soft RPV



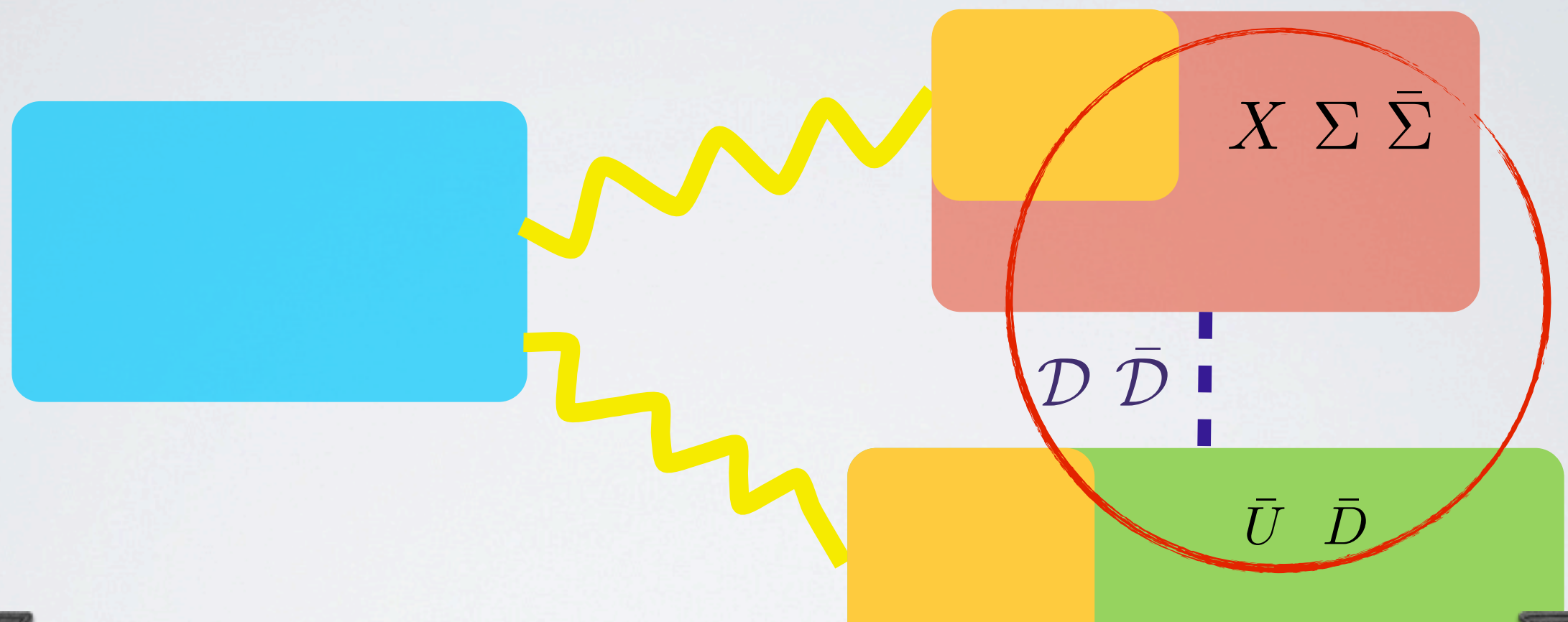
Breaking & mediation

soft RPV



Breaking & mediation

soft RPV



the potential $W \supset \bar{U}\bar{D}\bar{D} + X\mathcal{D}\bar{D} + X\Sigma\bar{\Sigma} + M_{\mathcal{D}}\bar{D}\mathcal{D}$
mediates the B-violation (induced by the R-symm breaking) to
the visible sector

The UDD A-term

	$SU(3)_c$	$U(1)_Y$	$U(1)_H$	R
\bar{U}	$\bar{3}$	$-2/3$	0	1
\bar{D}	$\bar{3}$	$1/3$	0	1
$\bar{\mathcal{D}}$	$\bar{3}$	$1/3$	0	0
\mathcal{D}	3	$-1/3$	0	2
X	1	0	0	-1
Σ	1	0	1	3/2
$\bar{\Sigma}$	1	0	-1	3/2

$$W \supset \bar{U}\bar{D}\bar{\mathcal{D}} + X\mathcal{D}\bar{D} + X\Sigma\bar{\Sigma} + M_{\mathcal{D}}\bar{\mathcal{D}}\mathcal{D}$$

integrating out the heavy mediator

$$W \supset \frac{\bar{U}\bar{D}\bar{\mathcal{D}}}{M_{\mathcal{D}}} X + X\Sigma\bar{\Sigma}$$

$\langle X \rangle = 0$, the RPV coupling only comes from the A-term

$$\mathcal{L}_{RPV} = \frac{\tilde{U}\tilde{D}\tilde{\mathcal{D}}}{M_{\mathcal{D}}} (\tilde{\Sigma}\tilde{\Sigma})^*$$

The UDD A-term

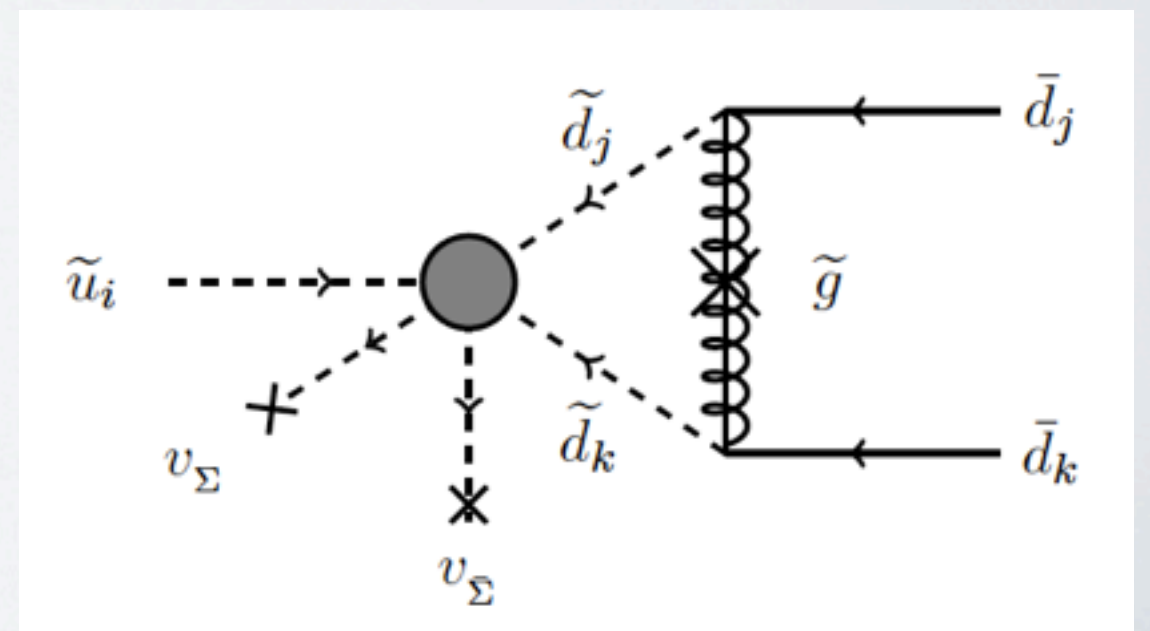
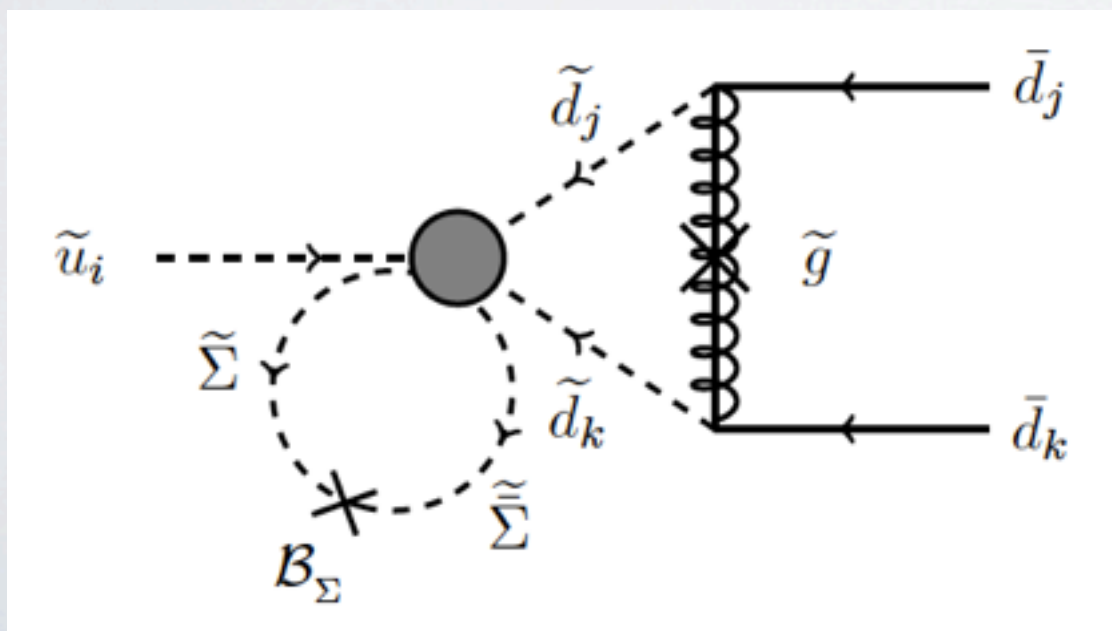
$$\mathcal{L}_{RPV} = \kappa_{i[j} \kappa'_{k]} \epsilon^{abc} \frac{\tilde{U}_a^i \tilde{D}_b^j \tilde{D}_c^k}{M_{\mathcal{D}}} (\tilde{\Sigma} \tilde{\Sigma})^*$$

R-breaking B-term

$$\mathcal{L} \subset \mathcal{B}_{\Sigma} \tilde{\Sigma} \tilde{\Sigma}$$

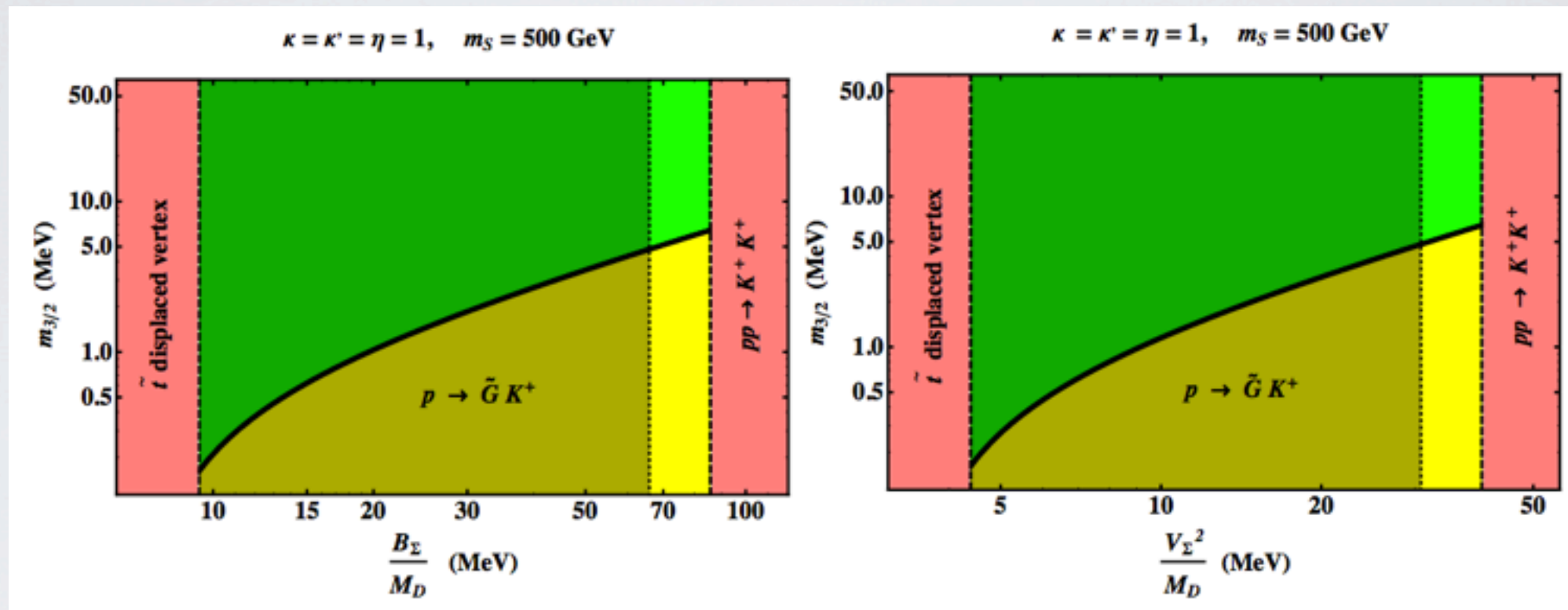
Radiatively induced VEV

$$\langle \tilde{\Sigma} \tilde{\Sigma} \rangle \simeq m_{soft}^2$$



Parameter space

For a generic flavor structure



For $\sqrt{B_\Sigma} \simeq v_\Sigma \simeq TeV$, the RPV mediator mass

$$M_D \simeq 10^{4-5} \text{ TeV}$$

$$M_D \sim \sqrt{F}$$

Gravitino Dark Matter

In models with baryonic RPV, the LSP can be stable if it's lighter than hadrons

- the $\mathcal{O}(10)$ MeV scale gravitino can be a thermally produced DM
- no additional non-thermal productions due to the RPV decay

$$\Omega_{3/2} h^2 \simeq 0.1 \left(\frac{T_R}{10^5 \text{ GeV}} \right) \left(\frac{m_{3/2}}{20 \text{ MeV}} \right)^{-1} \left(\frac{M_{\tilde{g}}}{800 \text{ GeV}} \right)^2$$

M.Boltz, A. Brandenburg and W. Buchmuller (2000)

To conclude

It's nice to think about the RPV

"softly"

In the example we show here

no additional symmetry required to forbid RPV (use R-symmetry)

no extra mass scales or small couplings put in by hand

gives a dark matter candidate

based on the gauge mediation setup

BACKUP

Experimental Constraints

upper bound on λ'' (want the process to happen slower)

di-nucleon decay $\tau_{pp \rightarrow KK} \geq 1.7 \times 10^{32} \text{ yrs}$

neutron/anti-neutron oscillation $\tau_{n-\bar{n}} \geq 2.44 \times 10^8 \text{ sec}$

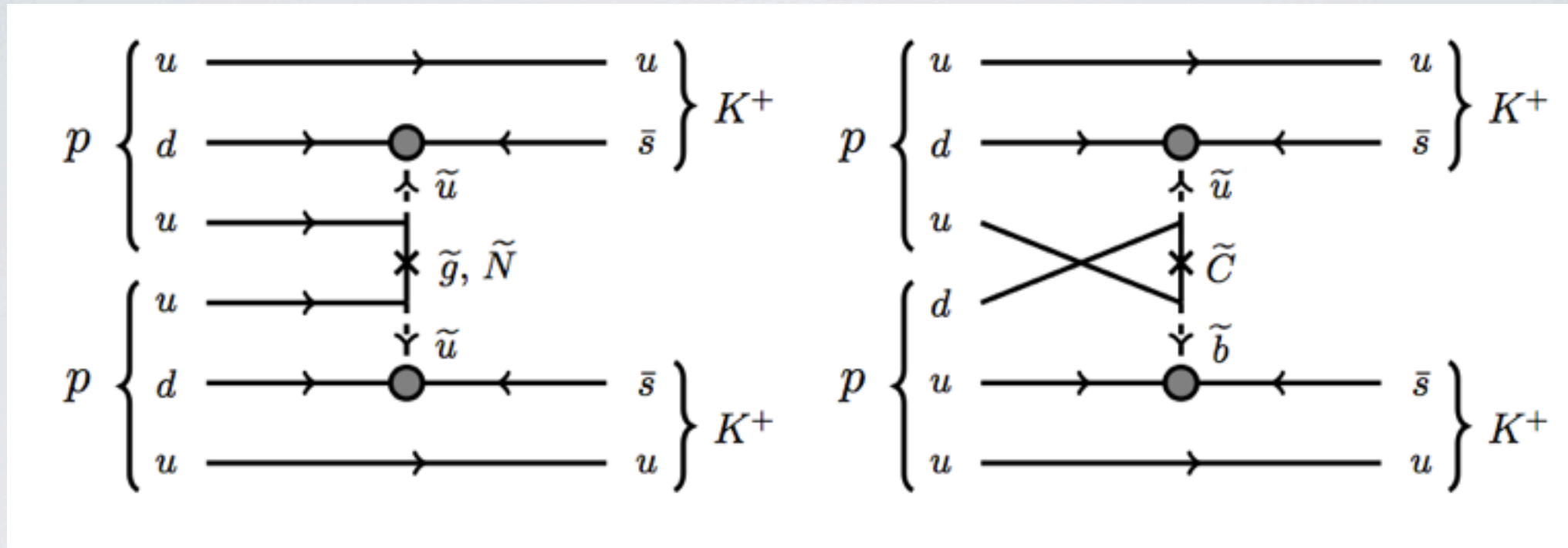
lower bound on λ'' (want the stop to decay faster)

stop decay length $l_{\tilde{t}} \leq 2 \text{ mm}$ (prompt) $l_{\tilde{t}} \leq 10 \text{ cm}$ (displaced)

lower bound on the ~~SUSY~~ scale (a smaller gravitino coupling)

proton decay into gravitino $\tau_{p \rightarrow K + \nu} \geq 2.3 \times 10^{33} \text{ yrs}$

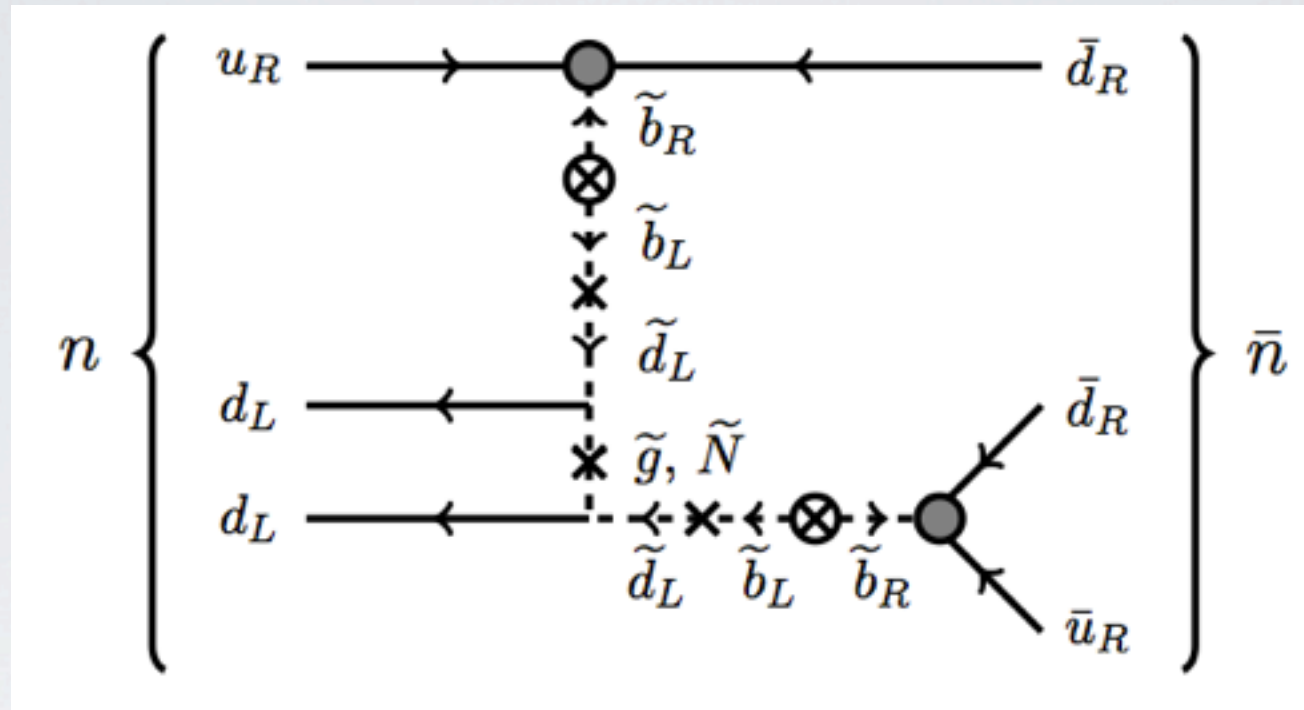
Di-nucleon decay



$$\Gamma_{pp \rightarrow KK} \sim \rho_N \frac{128 \pi \alpha_s^2 \Lambda^{10}}{m_p^2 m_{\tilde{u}}^8 M_{\tilde{g}}^2} (\lambda''_{uds})^2$$

$$\lambda''_{uds} \lesssim 2.5 \times 10^{-7} \left(\frac{150 \text{ MeV}}{\Lambda} \right)^{5/2} \left(\frac{M_{\tilde{g}}}{800 \text{ GeV}} \right)^{1/2} \left(\frac{m_{\tilde{u}}}{500 \text{ GeV}} \right)^2$$

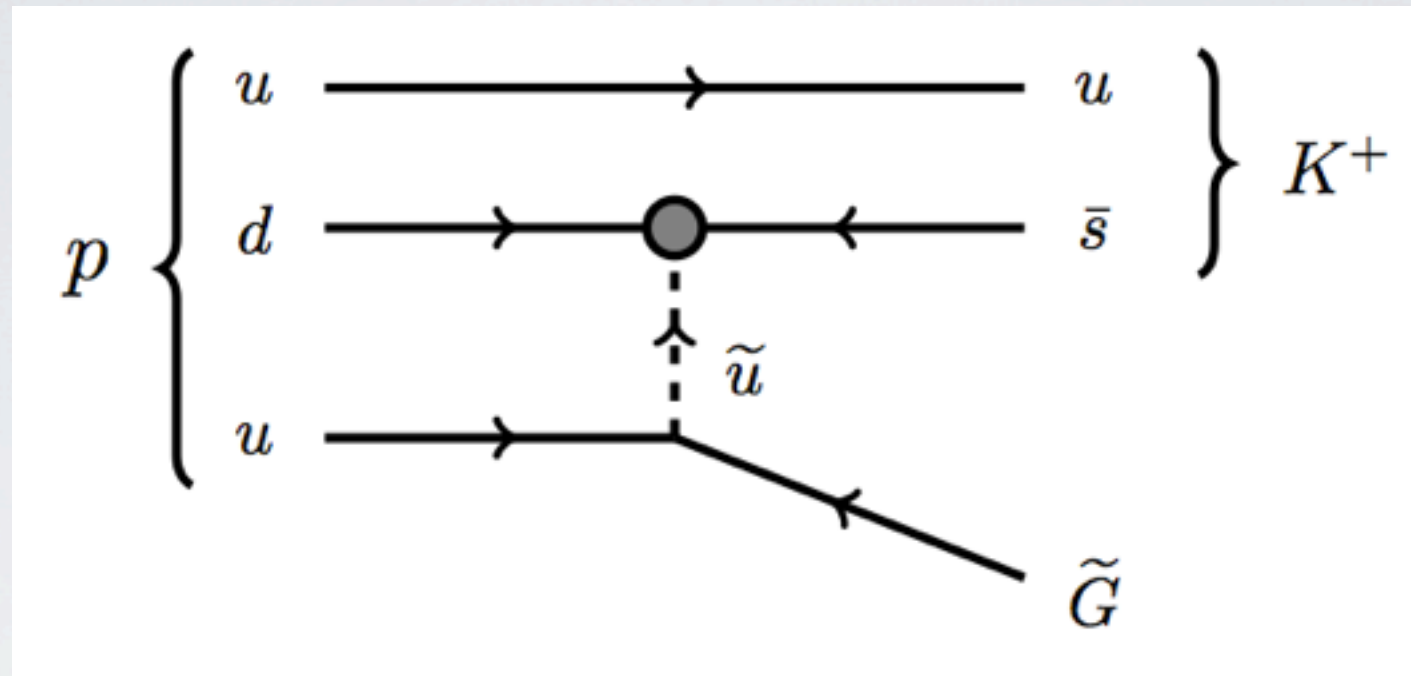
Neutron/anti-neutron oscillation



$$\mathcal{M}_{n-\bar{n}} \sim g_s^2 \epsilon^2 \lambda^6 \Lambda \left(\frac{\Lambda}{m_{\tilde{q}}} \right)^4 \left(\frac{\Lambda}{M_{\tilde{g}}} \right) (\lambda''_{udb})^2$$

$$\lambda''_{udb} \lesssim 1.7 \times 10^{-6} \epsilon^{-2} \left(\frac{m_{\tilde{q}}}{500 \text{ GeV}} \right)^4 \left(\frac{250 \text{ MeV}}{\Lambda} \right)^6 \left(\frac{M_{\tilde{g}}}{800 \text{ GeV}} \right)$$

Proton decay

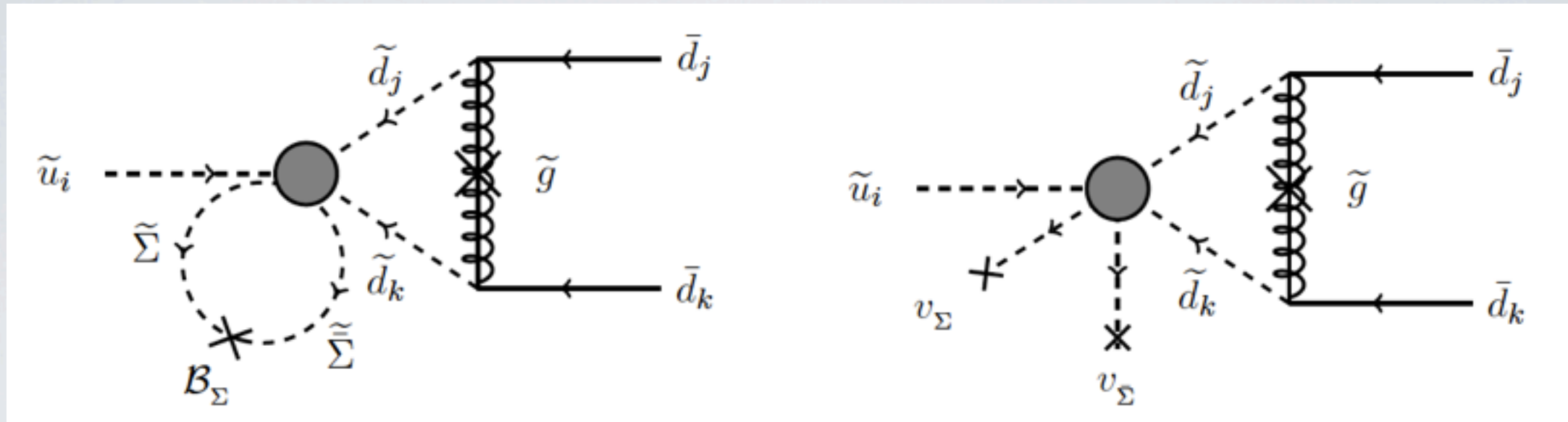


$$\Gamma_{p \rightarrow K + \tilde{G}} \sim \frac{m_p}{8\pi} \left(\frac{\Lambda}{m_{\tilde{u}}} \right)^4 \left(\frac{\Lambda^2}{\sqrt{3} m_{3/2} M_{pl}} \right)^2 (\lambda''_{uds})^2$$

$$m_{3/2} \geq 4.7 \text{ MeV} \left(\frac{\Lambda}{250 \text{ MeV}} \right)^4 \left(\frac{500 \text{ GeV}}{m_{\tilde{u}}} \right)^2 \left(\frac{\lambda''_{uds}}{10^{-7}} \right)$$

$$\sqrt{F} \gtrsim 3.2 \times 10^5 \text{ TeV}$$

The stop decay length



$$\Gamma_{\tilde{t} \rightarrow q\bar{q}} = \frac{m_{\tilde{t}}}{8\pi} \sin^2 \theta_{\tilde{t}} |\lambda''_{tqq}|^2$$

displaced jets are hard to see

$$\lambda''_{tds} > (0.26 - 1.8) \times 10^{-7} \left(\frac{300 \text{ GeV}}{m_{\tilde{t}}} \right)^{1/2}$$

bound for $\ell_{\tilde{t}} \leq 10 \text{ cm}$

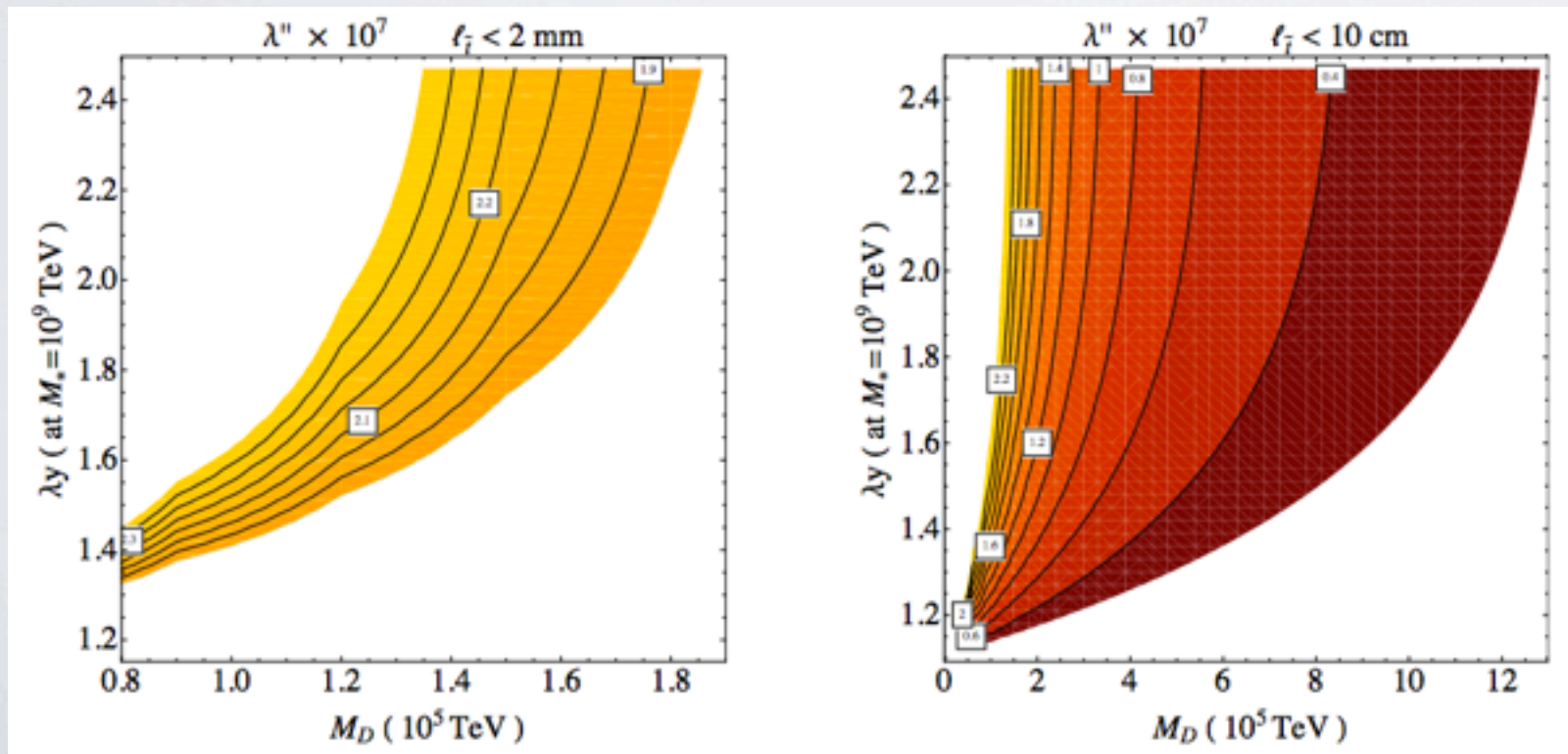
bound for $\ell_{\tilde{t}} \leq 2 \text{ mm}$

Mediator mass in the $\langle \tilde{\Sigma} \rangle$ model

to obtain $\langle \tilde{\Sigma} \rangle$, we can generate the tachyonic mass $-m_{\Sigma}^2$ through the RG running with the help of extra matter couplings. for example,

$$W \supset \eta \Sigma X \bar{\Sigma} + \lambda_Y \Sigma Y^2 + \lambda_{\bar{Y}} \bar{\Sigma} \bar{Y}^2$$

The size of $\langle \tilde{\Sigma} \rangle$ + experimental bounds set the upper/lower bounds on the mediator mass



here we assume

$$\sqrt{F} = 4 \times 10^5 \text{ TeV}$$

from these plot

$$M_D \sim \sqrt{F}$$

is allowed!