

**SUSY 2013**

21st International Conference on Supersymmetry and Unification of Fundamental Interactions

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# Probing Supersymmetry in B meson decays facing the recent LHC data

26 August 2013

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Y.Shimizu, M.Tanimoto, K.Y. [arXiv: 1307.0374]

# Introduction

The LHCb collaboration has reported new data of the CP violation of Bs meson.

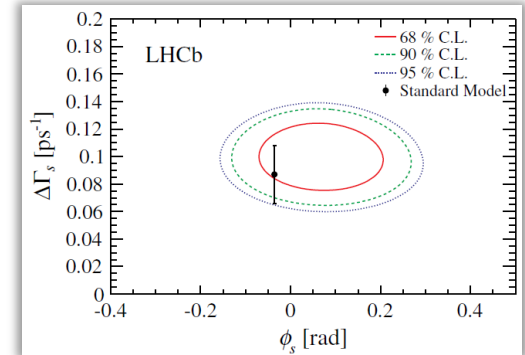
Time dependent CP asymmetry in  $B_s^0 \rightarrow J/\psi K^+ K^-$

[LHCb, Phys. Rev. D 87, 112010(2013)]

$$\phi_s(\text{Exp}) = 0.07 \pm 0.09 \pm 0.01$$

$$\phi_s(\text{SM}) = -0.0363 \pm 0.0017 \quad [\text{The CKMfitter, 2011}]$$

$$\text{Asymmetry} = \sin \phi_s$$



Time dependent CP asymmetry in  $B_s \rightarrow \phi\phi$

[LHCb, Phys.Rev.Lett.110, 241802 (2013)]

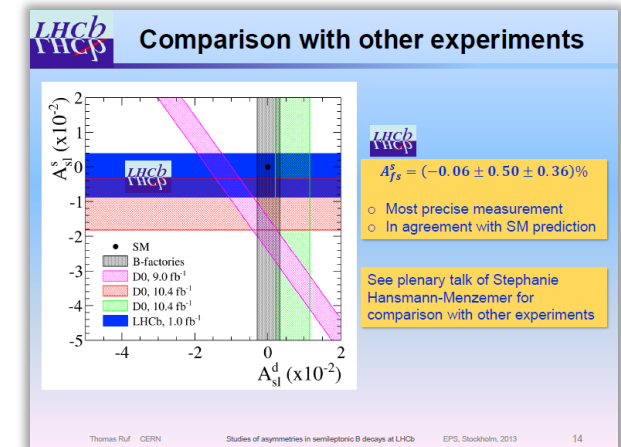
$$\phi_s^{\phi\phi}(\text{Exp}) = [-2.46, -0.78]\text{rad} \quad (68\% \text{C.L.})$$

$$\phi_s^{\phi\phi}(\text{SM}) \sim \mathcal{O}(\lambda^4)$$

Semi-leptonic asymmetry  $a_{sl}^s$  [EPS conference 2013]

$$a_{sl}^s(\text{Exp}) = (-0.06 \pm 0.50 \pm 0.36)\%$$

$$a_{sl}^s(\text{SM}) = (0.0019 \pm 0.0003)\% \quad [\text{A.Lenz, 2012}]$$



# Introduction

## Searching for SUSY particle at LHC

The SUSY signals have not been observed yet.

The lower bounds for the susy-particle masses have pushed up to TeV scale.

95% CL exclusions for (best) massless LSP scenarios:

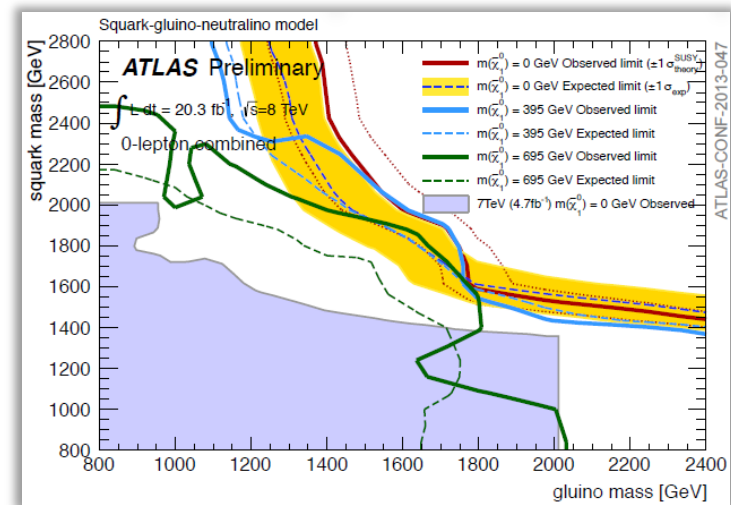
$$m(\tilde{g}) < 1300 \text{ GeV}$$

$$m(\tilde{q}) < 1400 \text{ GeV}$$

$$m(\tilde{b}) < 650 \text{ GeV}$$

$$m(\tilde{t}) < 680 \text{ GeV}$$

$$m(\tilde{\ell}_L) < 300 \text{ GeV}$$



[Lepton-Photon symposium 2013 06]

Indirect search for SUSY is important.

**B physics can become the indirect search for SUSY.**  
**We discuss the sensitivity of the SUSY contribution to B decays.**

# Setup

**Natural SUSY scenario** [A.G.Cohen, D.B.Kaplan and A.E.Nelson, PLB 388 (1996) 588]

The first and second family squarks are very heavy, at  $O(10)$  TeV, on the other hand, the third family squark masses are at  $O(1)$  TeV.

$$m_{\tilde{q}_{1,2}} = \mathcal{O}(10) \text{ TeV} \quad m_{\tilde{b}_1} = 1 \text{ TeV}, m_{\tilde{b}_2} = 1.1 \text{ TeV} \quad m_{\tilde{g}} = 2 \text{ TeV}$$

## The gluino-squark-quark interaction

$$\mathcal{L}_{\text{int}}(\tilde{g}q\tilde{q}) = -i\sqrt{2}g_s \sum_{\{q\}} \tilde{q}_i^* (T^a) \overline{\tilde{G}}^a \left[ (\Gamma_{GL}^{(q)})_{ij} \mathbf{L} + (\Gamma_{GR}^{(q)})_{ij} \mathbf{R} \right] q_j + \text{h.c.}$$

**Mixing matrix** We work in the basis of mass eigenstate.

$$\Gamma_{GL}^{(d)} = \begin{pmatrix} 1 & 0 & \delta_{13}^{dL} c_\theta & 0 & 0 & -\delta_{13}^{dL} s_\theta e^{i\phi} \\ 0 & 1 & \delta_{23}^{dL} c_\theta & 0 & 0 & -\delta_{23}^{dL} s_\theta e^{i\phi} \\ -\delta_{13}^{dL*} & -\delta_{23}^{dL*} & c_\theta & 0 & 0 & -s_\theta e^{i\phi} \end{pmatrix} \begin{pmatrix} d_L \\ s_L \\ b_L \end{pmatrix}$$

$$\Gamma_{GR}^{(d)} = \begin{pmatrix} 0 & 0 & \delta_{13}^{dR} s_\theta e^{-i\phi} & 1 & 0 & \delta_{13}^{dR} c_\theta \\ 0 & 0 & \delta_{23}^{dR} s_\theta e^{-i\phi} & 0 & 1 & \delta_{23}^{dR} c_\theta \\ 0 & 0 & s_\theta e^{-i\phi} & -\delta_{13}^{dR*} & -\delta_{23}^{dR*} & c_\theta \end{pmatrix} \begin{pmatrix} d_R \\ s_R \\ b_R \end{pmatrix}$$

$$\begin{matrix} \tilde{d}_1 & \tilde{s}_1 & \tilde{b}_1 & \tilde{d}_2 & \tilde{s}_2 & \tilde{b}_2 \end{matrix}$$

We assume the mixing parameters  $\delta_{13}^{dL}, \delta_{23}^{dL}, \delta_{13}^{dR}, \delta_{23}^{dR}$  as  $|\delta_{13}^{dL}| = |\delta_{13}^{dR}|, |\delta_{23}^{dL}| = |\delta_{23}^{dR}|$ .

We assume the mixing angle  $\theta = 10^\circ \sim 35^\circ$ .

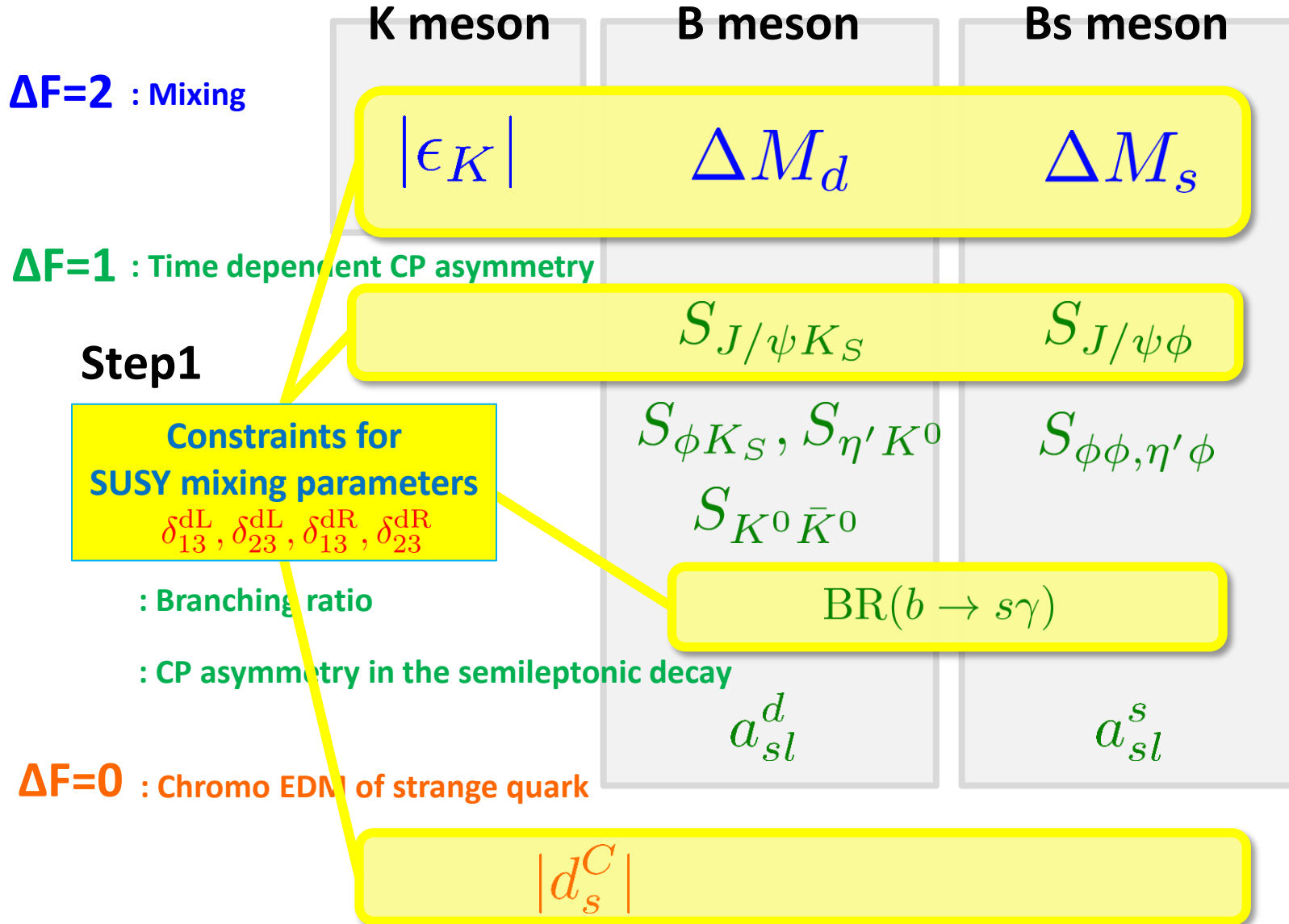
(Mixing angle between the left-handed sbottom and the right-handed one.)

# Targets

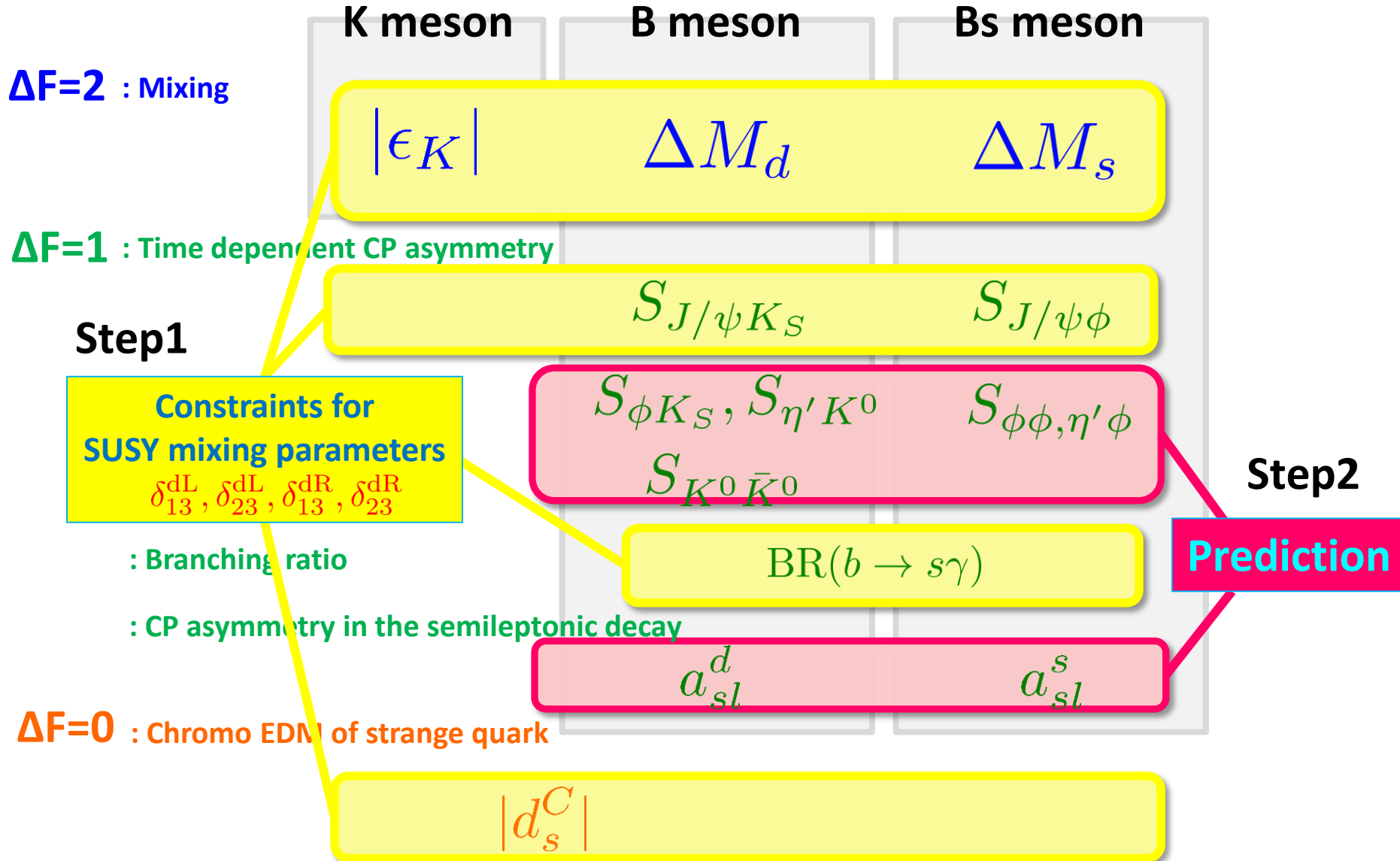
	K meson	B meson	Bs meson
$\Delta F=2$ : Mixing	$ \epsilon_K $	$\Delta M_d$	$\Delta M_s$
$\Delta F=1$ : Time dependent CP asymmetry		$S_{J/\psi K_S}$	$S_{J/\psi \phi}$
		$S_{\phi K_S}, S_{\eta' K^0}$	$S_{\phi\phi, \eta' \phi}$
		$S_{K^0 \bar{K}^0}$	
: Branching ratio		BR( $b \rightarrow s\gamma$ )	
: CP asymmetry in the semileptonic decay		$a_{sl}^d$	$a_{sl}^s$
$\Delta F=0$ : Chromo EDM of strange quark			

$$|d_s^C|$$

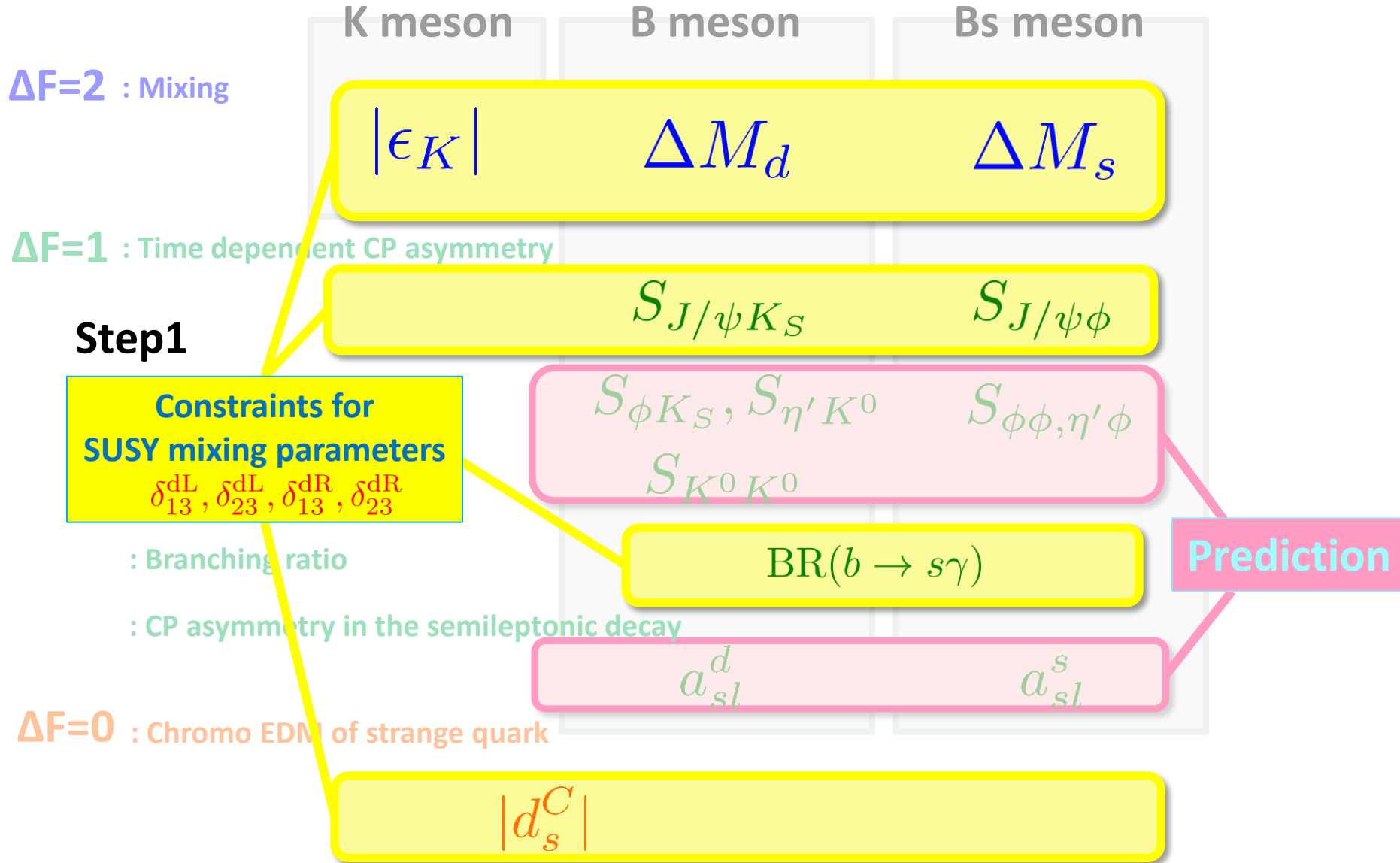
# Targets



# Targets

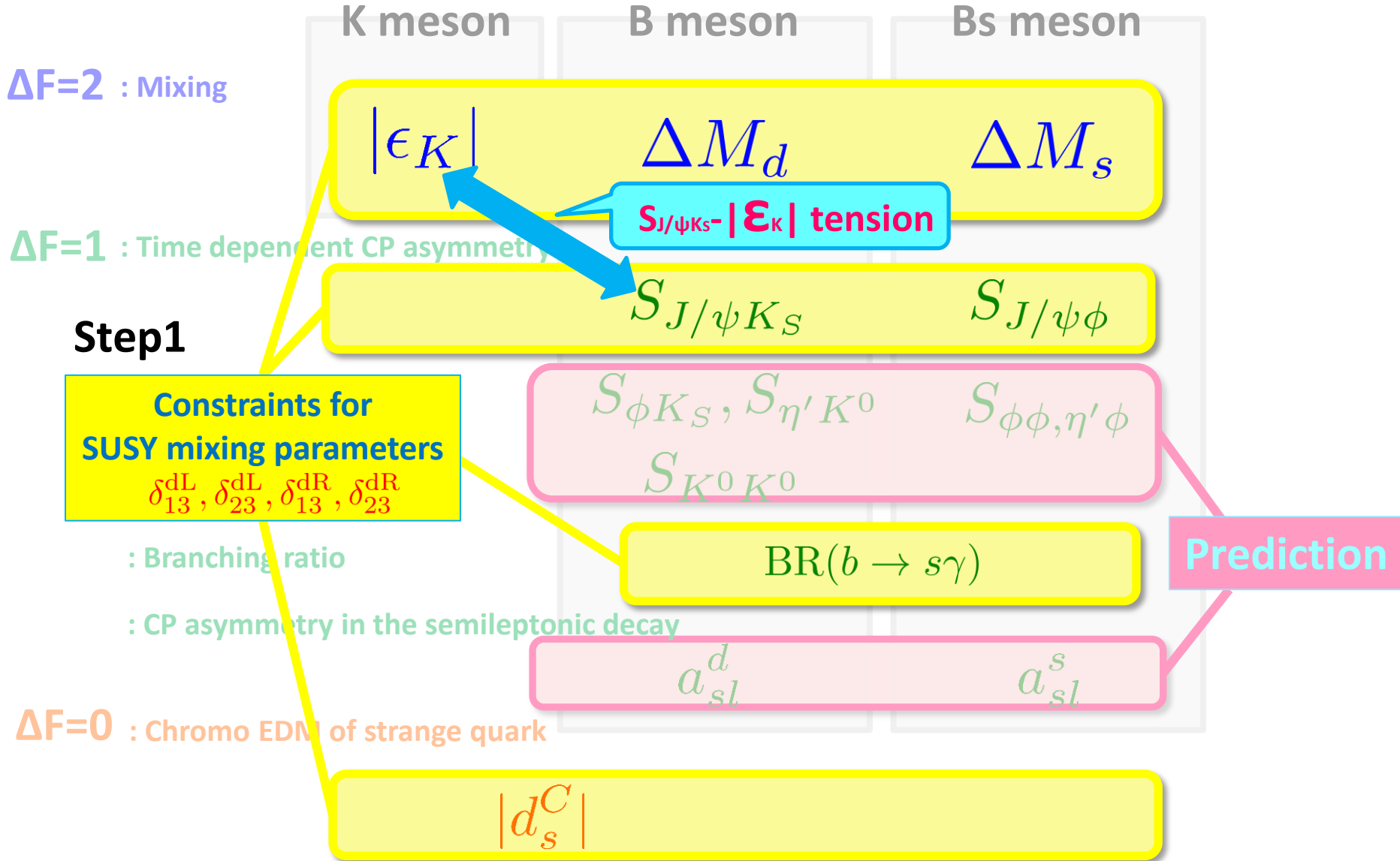


# Targets





# Targets



# $S_{J/\psi K_s} - |\epsilon_K|$ tension

[Andrzej J. Buras and Diego Guadagnoli, 2008]  
[Andrzej J. Buras, 2011]

$|\epsilon_K|$  is given in terms of  $\sin(2\beta)$  because there is only one CP violating phase in the SM.

$$\text{SM} \quad |\epsilon_K| \propto \hat{B}_K |V_{cb}|^4 \sin(2\beta_{\text{SM}}) \quad S_{J/\psi K_s} = \sin(2\beta_{\text{exp}}) = \sin(2\beta_{\text{SM}})$$



**B<sub>K</sub> parameter**

$$B_K = \frac{\langle \bar{K}^0 | O^{\Delta S=2} | K^0 \rangle}{\frac{8}{3} \langle \bar{K}^0 | A_\mu | 0 \rangle \langle 0 | A^\mu | K^0 \rangle} \quad \hat{B}_K \equiv [\alpha_s(\mu)]^{-2/9} \left[ 1 + \frac{\alpha_s(\mu)}{4\pi} J_3 \right] B_K(\mu)$$

**Recent lattice work** [C.Aubin, et all, 2009]

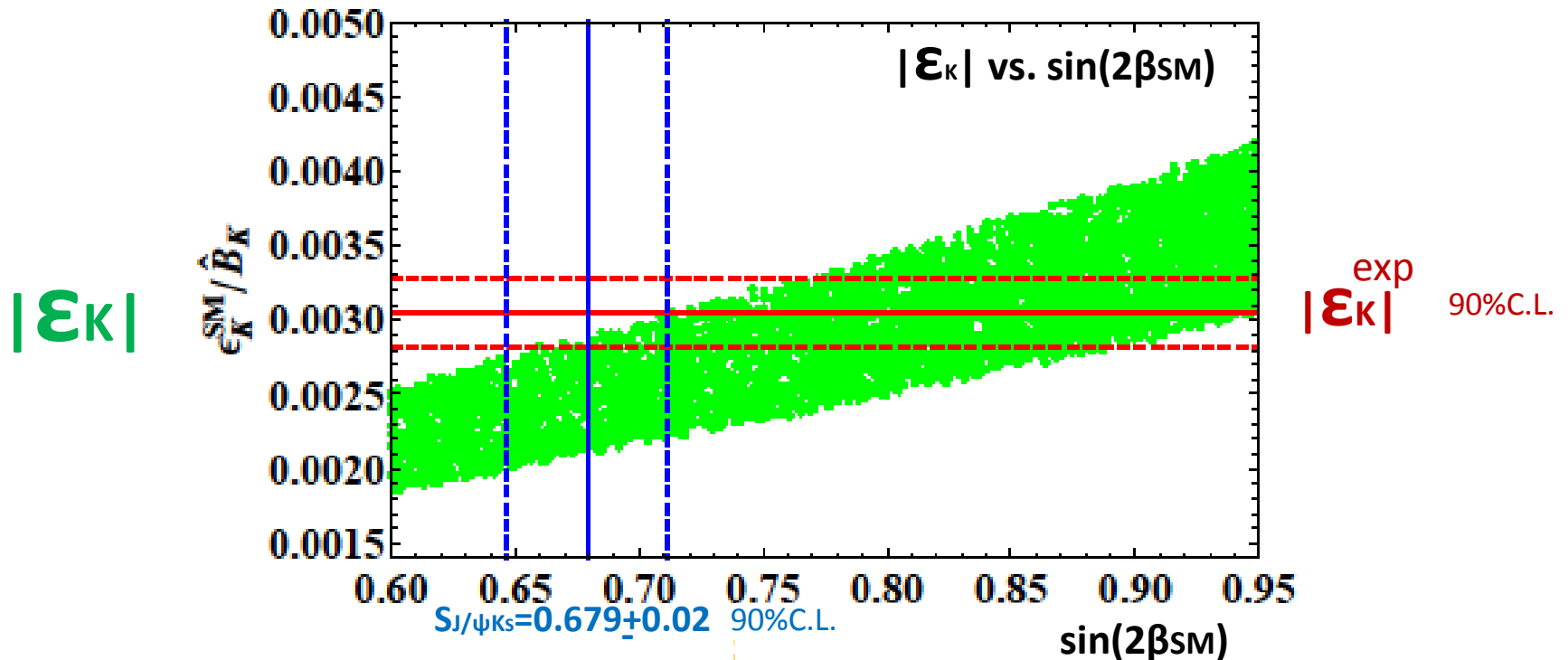
$$\hat{B}_K = 0.73 \pm 0.03$$

# $S_{J/\psi K_s} - |\epsilon_K|$ tension

[Andrzej J. Buras and Diego Guadagnoli, 2008]  
[Andrzej J. Buras, 2011]

$|\epsilon_K|$  is given in terms of  $\sin(2\beta)$  because there is only one CP violating phase in the SM.

**SM**  $|\epsilon_K| \propto \hat{B}_K |V_{cb}|^4 \sin(2\beta_{SM})$   $S_{J/\psi K_s} = \sin(2\beta_{exp}) = \sin(2\beta_{SM})$

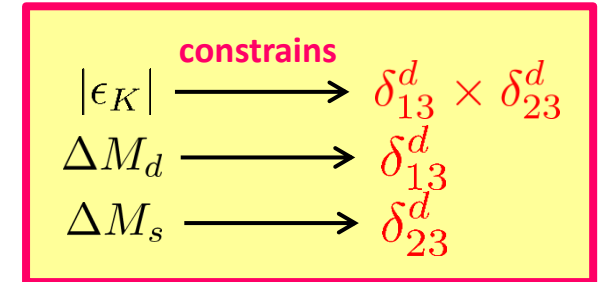
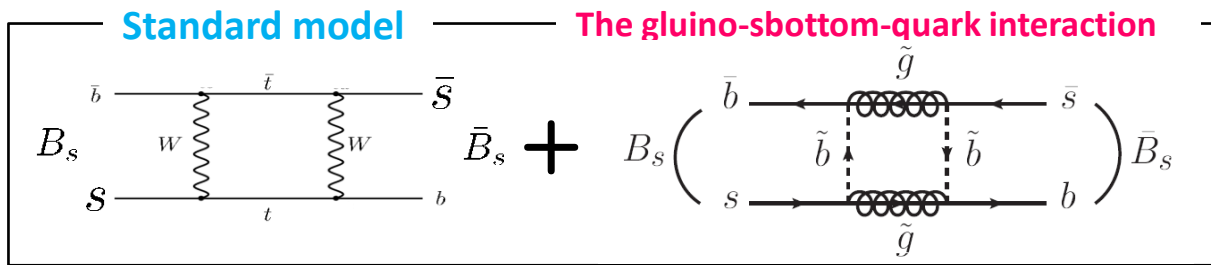


It is noticed that the consistency between the SM prediction and the experimental data in  $\sin(2\beta)$  and  $|\epsilon_K^{SM}| / \hat{B}_K$  is marginal.

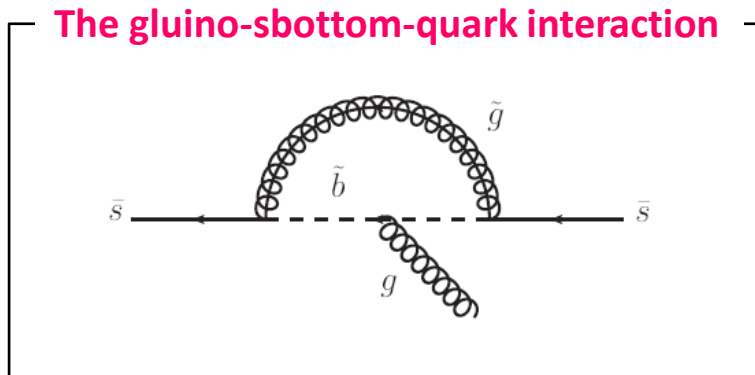
We will show that this tension is understood by taking account of the SUSY box diagram through **the gluino-sbottom-quark interaction**.

# Constraints

$\Delta F=2$  process : K- $\bar{K}$ , B- $\bar{B}$ ,  $B_s$ - $\bar{B}_s$  mixing  $|\epsilon_K| \Delta M_d \Delta M_s$

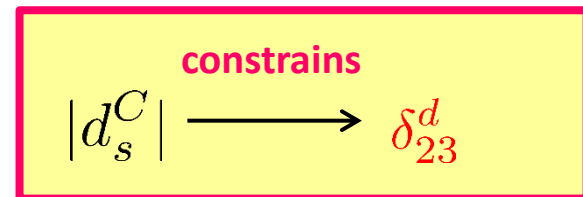


$\Delta F=0$  process : Chromo EDM of strange quark  $|d_s^C|$



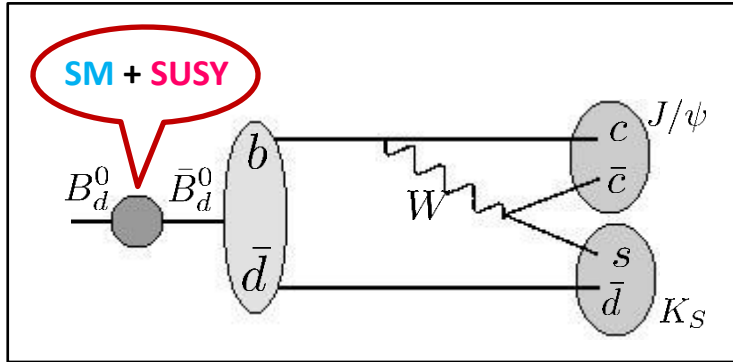
[K.Fuyuto, J.Hisano and N.Nagata, Phys.Rev.D 87 (2013)]

$$e|d_s^C| < 0.5 \times 10^{-25} \text{ ecm}$$



# Constraints

**$\Delta F=1$  process** : Time dependent CP asymmetry  $S_{J/\psi K_S}$   $S_{J/\psi \phi}$



The dispersive part of mixing

$$\begin{aligned} M_{12}^q &= M_{12}^{q,\text{SM}} + M_{12}^{q,\text{SUSY}} \\ &= M_{12}^{q,\text{SM}} (1 + h_q e^{2i\sigma_q}) \\ &\quad (q = d, s) \end{aligned}$$

$$S_f = \frac{2\text{Im}(\lambda_f)}{1 + |\lambda_f|^2} \left[ \lambda = \frac{q}{p} \frac{\bar{A}^{\text{SM}} + \bar{A}^{\text{SUSY}}}{A^{\text{SM}} + A^{\text{SUSY}}} \right]$$

$$\frac{q}{p} \simeq \sqrt{\frac{M_{12}^*}{M_{12}}} = \sqrt{\frac{M_{12}^{\text{SM}*}}{M_{12}^{\text{SM}}}} \sqrt{\frac{1 + h_d e^{-2i\sigma_d}}{1 + h_d e^{2i\sigma_d}}}$$

SM

SM + SUSY

$$S_{J/\psi K_S} = \sin(2\beta_{\text{SM}}) \longrightarrow \sin(2\beta_{\text{SM}} + \text{Arg}(1 + h_d e^{2i\sigma_d}))$$

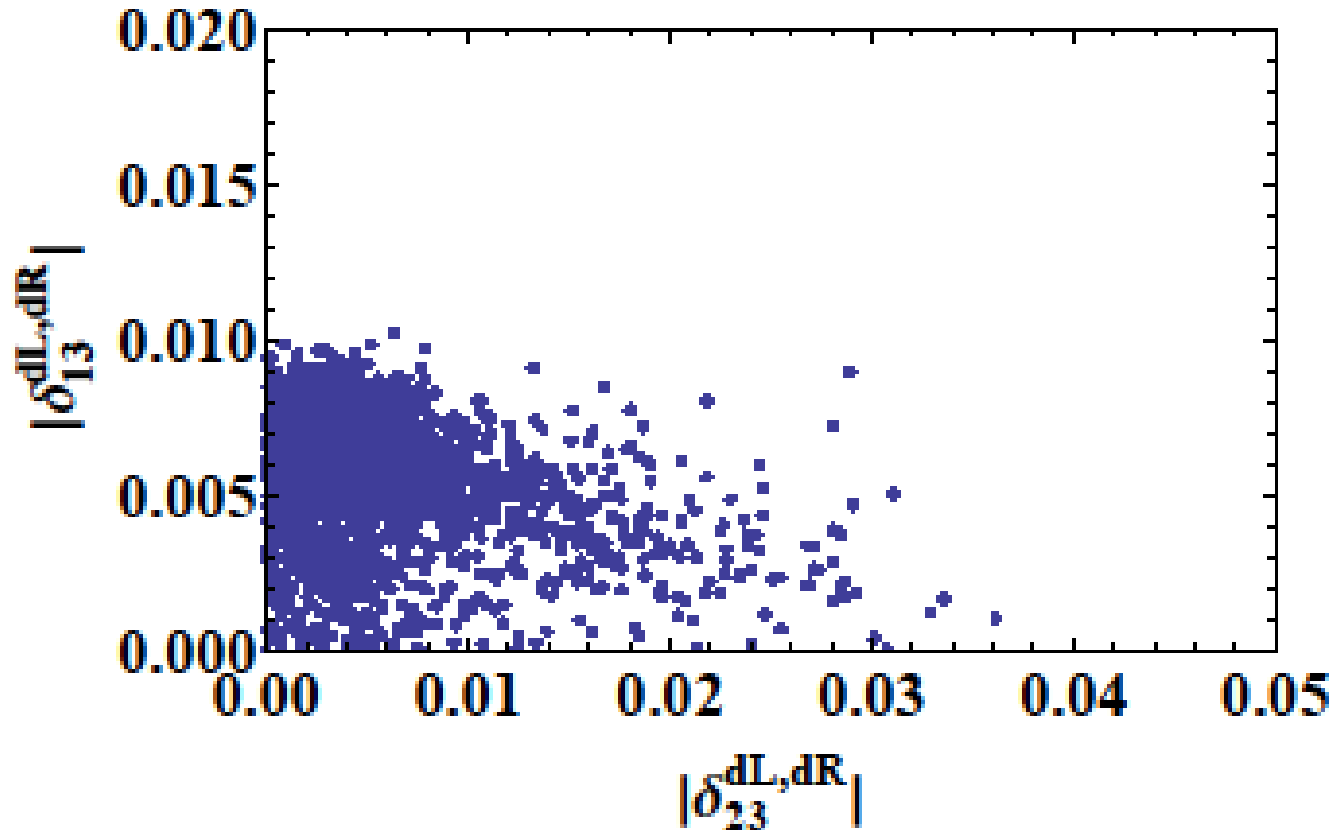
$$S_{J/\psi \phi} = \sin(-2\beta_{s\text{SM}}) \longrightarrow \sin(-2\beta_{s\text{SM}} + \text{Arg}(1 + h_s e^{2i\sigma_s}))$$

constrains

$$\begin{aligned} S_{J/\psi K_S} &\longrightarrow \delta_{13}^d \\ S_{J/\psi \phi} &\longrightarrow \delta_{23}^d \end{aligned}$$

# Numerical results

Allowed region of  $|\delta_{13}|, |\delta_{23}|$

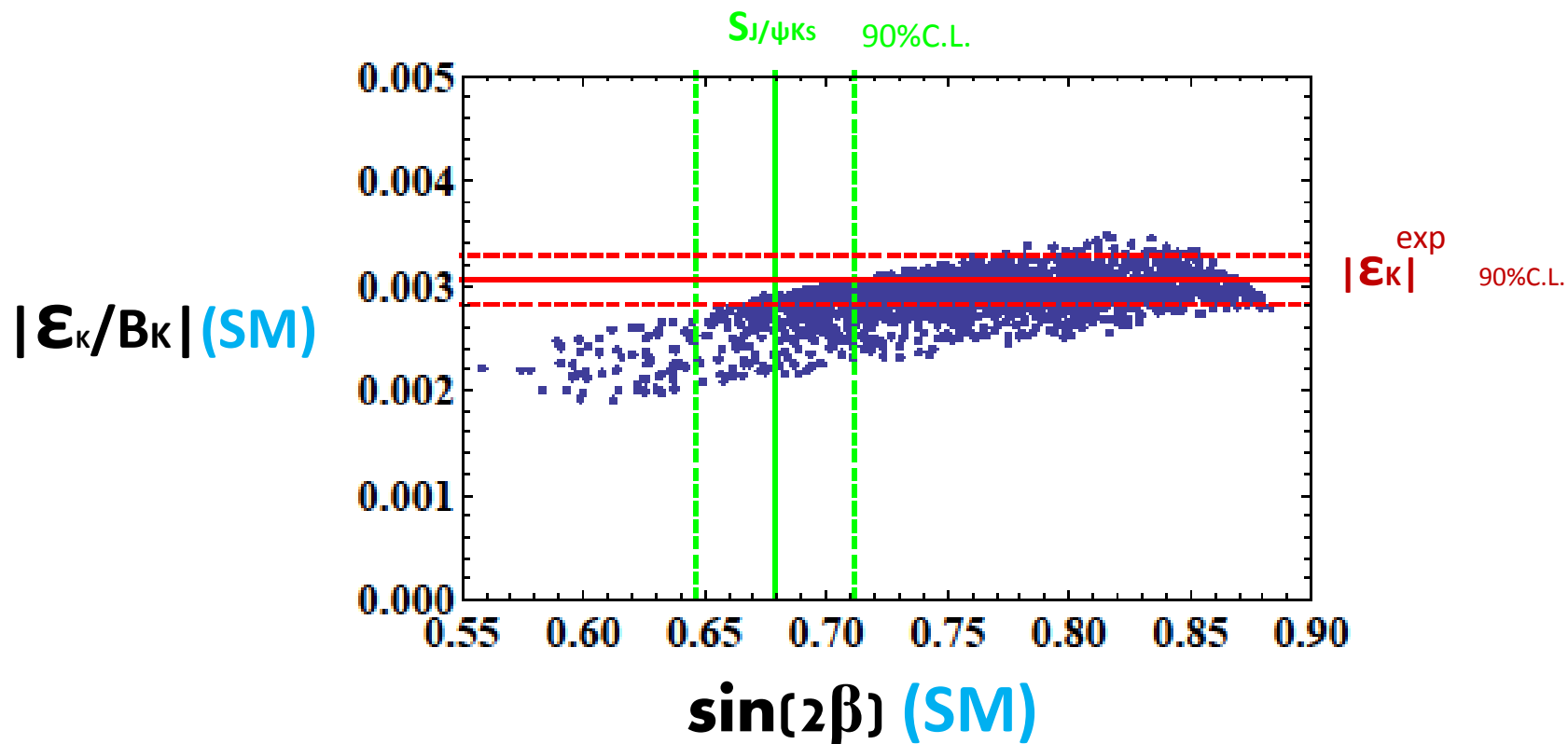


$$|\delta_{13}^{dL(dR)}| = 0 \sim 0.01 \quad |\delta_{23}^{dL(dR)}| = 0 \sim 0.04$$

@  $m_{\tilde{b}_1} = 1 \text{ TeV}, m_{\tilde{b}_2} = 1.1 \text{ TeV} \quad m_{\tilde{g}} = 2 \text{ TeV} \quad \theta = 10^\circ \sim 35^\circ$

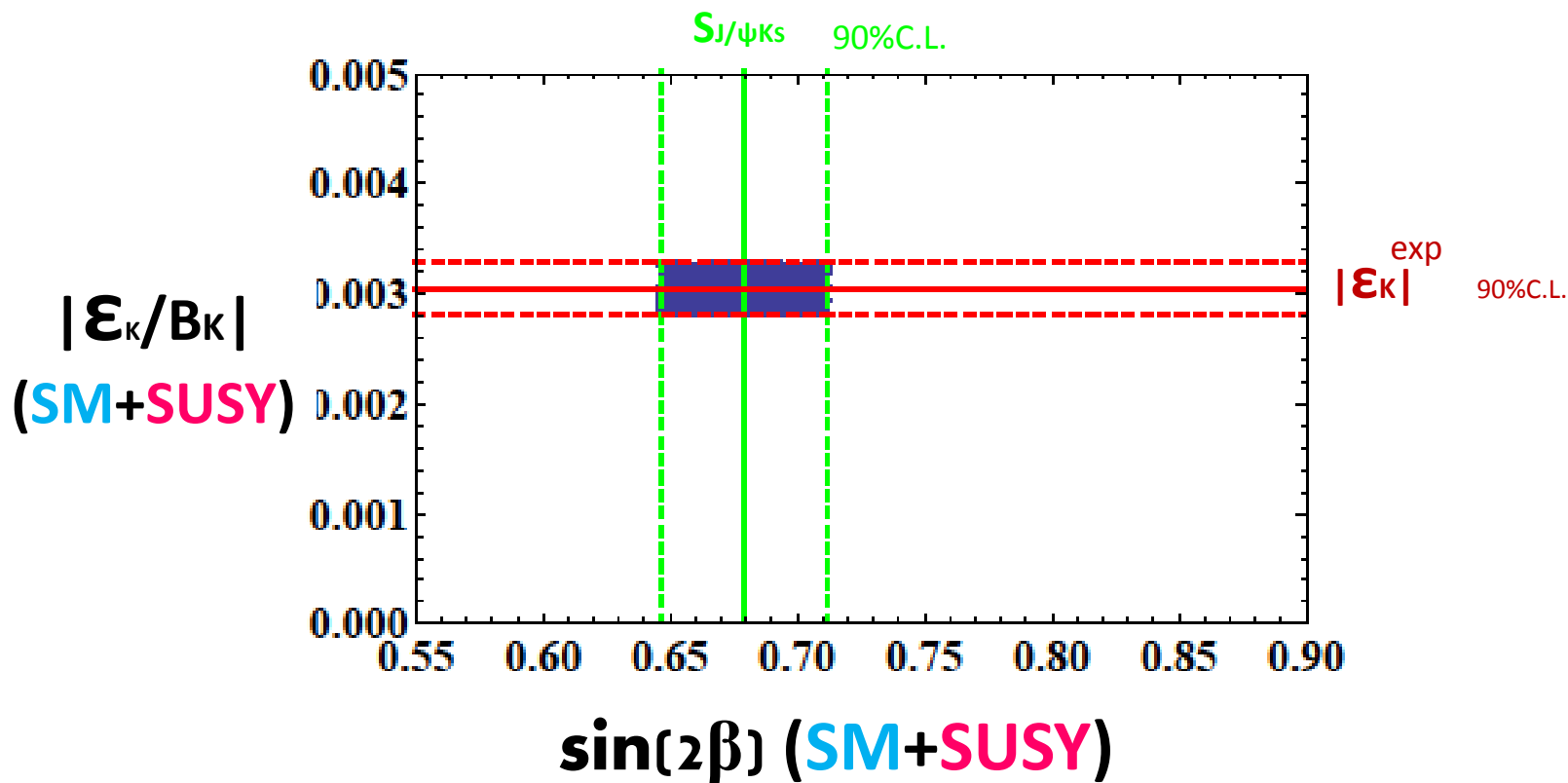
# Numerical results

$S_{J/\psi K_S} - |\epsilon_K|$  tension



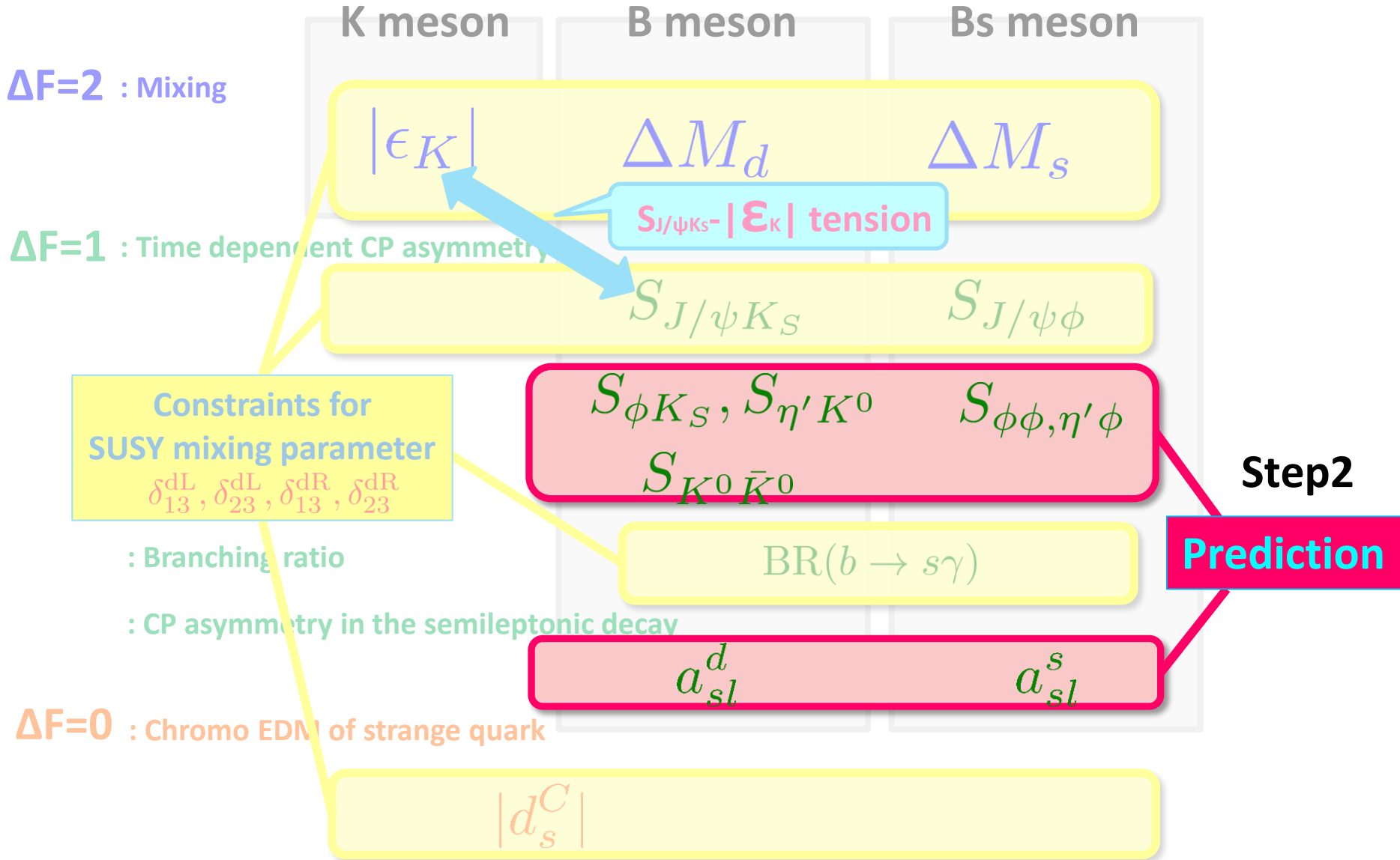
# Numerical results

$S_{J/\psi K_S} - |\epsilon_K|$  tension





# Targets



$S_f$ Time dependent CP asymmetry in  $B_q \rightarrow f$ 

$$\begin{aligned}
 A &= \frac{\Gamma(B^0(t) \rightarrow f) - \Gamma(\bar{B}^0(t) \rightarrow f)}{\Gamma(B^0(t) \rightarrow f) + \Gamma(\bar{B}^0(t) \rightarrow f)} \\
 &= \frac{|\lambda|^2 - 1}{|\lambda|^2 + 1} \cos(\Delta m_B t) + \frac{2\text{Im}\lambda}{|\lambda|^2 + 1} \sin(\Delta m_B t)
 \end{aligned}$$

$$\lambda = \frac{q}{p} \frac{\bar{A}^{SM} + \bar{A}^{SUSY}}{A^{SM} + A^{SUSY}}$$

mixing part     amplitude part

$$A(B_q \rightarrow f) \simeq \sum_i C_i O_i$$

$C_i$  : Wilson coefficients

$$C_i = C_i(\text{SM}) + C_i(\text{NP})$$

$O_i$  : Local operator

$$\lambda_{\phi K_S, \eta' K^0} = \text{Sin}(2\beta_{SM} + \text{Arg}(1 + h_d e^{2i\sigma_d})) \frac{\sum_{i=3-6,7\gamma,8G} (C_i^{\text{SM}} \langle O_i \rangle + C_i^{\tilde{g}} \langle O_i \rangle + \tilde{C}_i^{\tilde{g}} \langle \tilde{O}_i \rangle)}{\sum_{i=3-6,7\gamma,8G} (C_i^{\text{SM}*} \langle O_i \rangle + C_i^{\tilde{g}*} \langle O_i \rangle + \tilde{C}_i^{\tilde{g}*} \langle \tilde{O}_i \rangle)}$$

We focus on a few modes which decay at **1-loop level in SM**, because it is expected that SUSY contribution will be enhanced.

# Predictions

$$S_{\phi K_S}, S_{\eta' K^0}$$

Time dependent CP asymmetry in  $B_d \rightarrow \phi K_S, \eta' K^0$

## SM prediction

$$S_{J/\psi K_S} \simeq S_{\phi K_S, \eta' K_S}$$

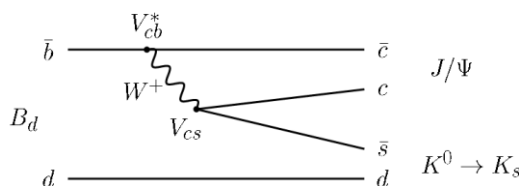
## Experimental results [HFAG,2012]

$$S_{J/\psi K_S} = 0.679 \pm 0.020$$

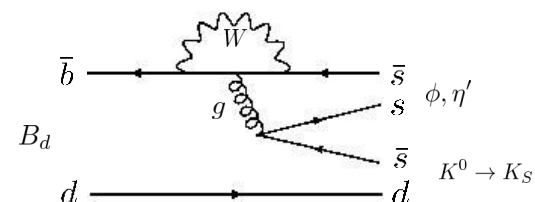
$$S_{\phi K_S} = 0.74^{+0.11}_{-0.13}$$

$$S_{\eta' K^0} = 0.59 \pm 0.07$$

$$B_d \rightarrow J/\psi K_S$$



$$B_d \rightarrow \phi K_S, \eta' K^0$$



Both CP violations come from CP phase in the  $B_d^0 - \bar{B}_d^0$  mixing.

## SUSY contribution

$$A^{SUSY}(\bar{B}_d \rightarrow \phi K_S) \propto C_{8G}^{\tilde{g}}(m_b) + \tilde{C}_{8G}^{\tilde{g}}(m_b)$$

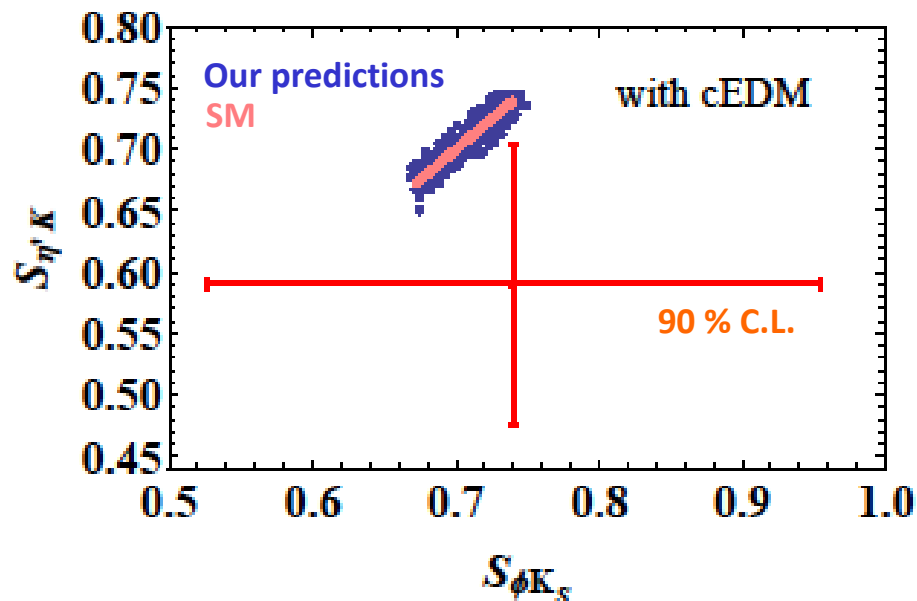
$$A^{SUSY}(\bar{B}_d \rightarrow \eta' K_S) \propto C_{8G}^{\tilde{g}}(m_b) - \tilde{C}_{8G}^{\tilde{g}}(m_b)$$

$$\tilde{C} : C(L \Leftrightarrow R) \quad \uparrow \text{Depends on } \delta_{23}^d$$

Difference of sign comes from parity of final state.

[M.Endo, S.Mishima and M.Yamaguchi, PLB 609 (2005)]

## Our prediction



# Predictions

$$S_{\phi\phi}, S_{\eta'\phi}$$

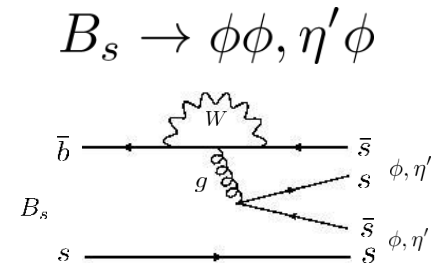
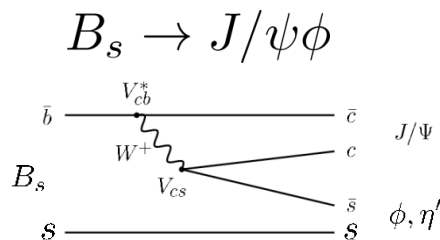
Time dependent CP asymmetry in  $B_s \rightarrow \phi\phi, \eta'\phi$

## SM prediction

$$S_{J/\psi\phi} = \sin\phi_s$$

$$\phi_s(\text{SM}) = -0.0363 \pm 0.0017$$

$$S_{\phi\phi, \eta'\phi} \simeq \mathcal{O}(\lambda^4)$$



## Experimental results

$$S_{J/\psi\phi} \quad [\text{LHCb, Phys. Rev. D 87, 112010(2013)}]$$

$$\phi_s(\text{Exp}) = 0.07 \pm 0.09 \pm 0.01$$

$$S_{\phi\phi} \quad [\text{LHCb, Phys.Rev.Lett.110, 241802 (2013)}]$$

$$\phi_s(\phi\phi) = [-2.46, -0.76] (68\% \text{C.L.}) \quad \text{New !}$$

## SUSY contribution

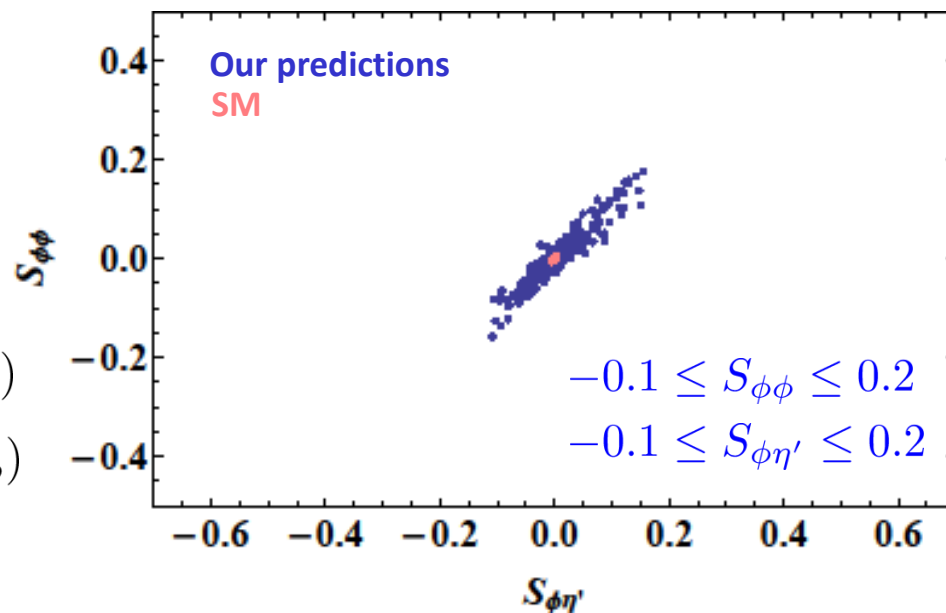
$$A^{SUSY}(\bar{B}_s \rightarrow \phi\phi) \propto C_{8G}^{\tilde{g}}(m_b) - \tilde{C}_{8G}^{\tilde{g}}(m_b)$$

$$A^{SUSY}(\bar{B}_s \rightarrow \phi\eta') \propto C_{8G}^{\tilde{g}}(m_b) + \tilde{C}_{8G}^{\tilde{g}}(m_b)$$



Depends on  $\delta_{23}^d$

## Our prediction



# Predictions

$$S_{\phi\phi}, S_{\eta'\phi}$$

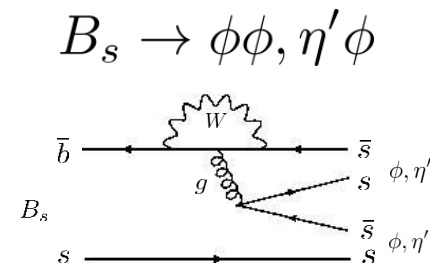
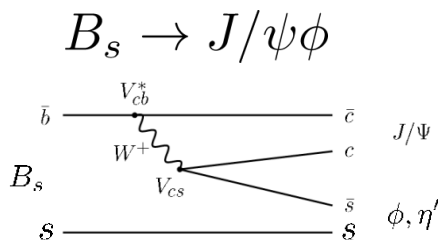
Time dependent CP asymmetry in  $B_s \rightarrow \phi\phi, \eta'\phi$

## SM prediction

$$S_{J/\psi\phi} = \sin \phi_s$$

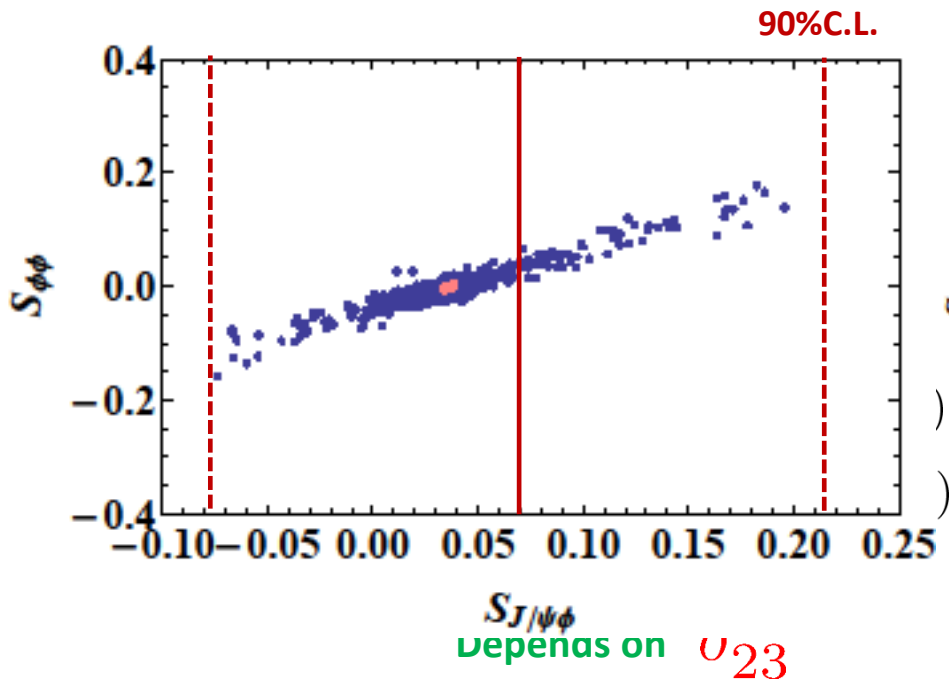
$$\phi_s(\text{SM}) = -0.0363 \pm 0.0017$$

$$S_{\phi\phi, \eta'\phi} \simeq \mathcal{O}(\lambda^4)$$

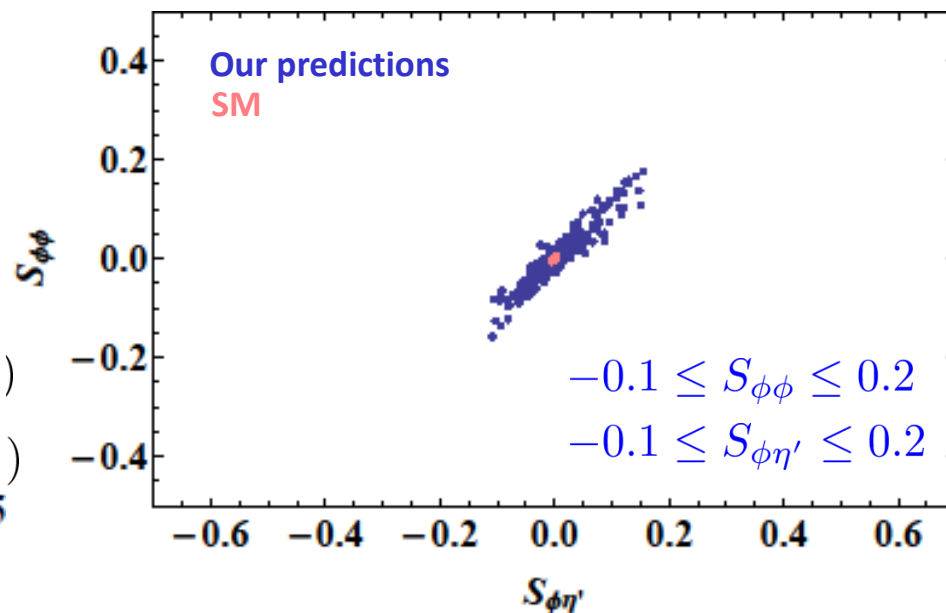


## Experimental results

$$S_{J/\psi\phi} \quad [\text{LHCb, Phys. Rev. D 87, 112010(2013)}]$$



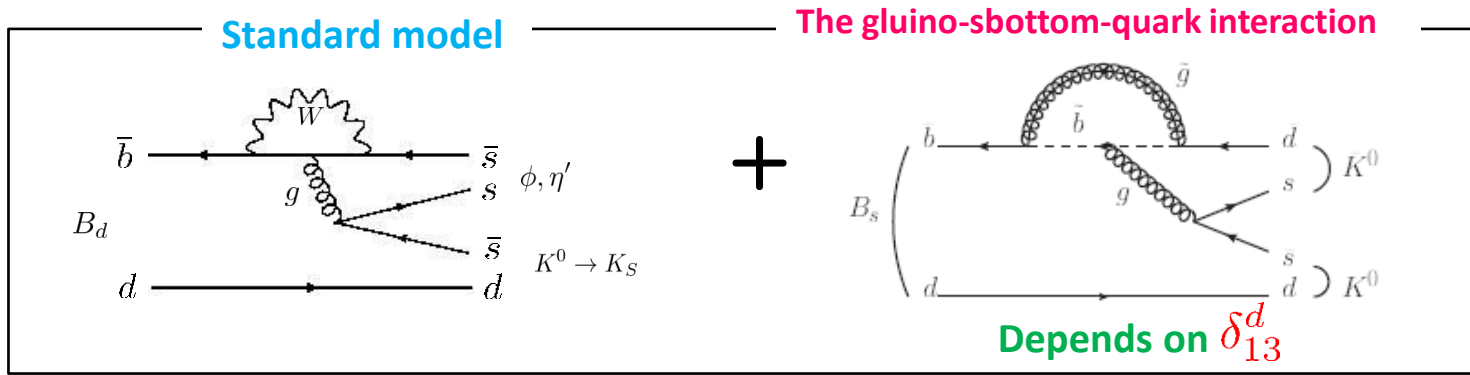
## Our prediction



# Predictions

$S_{K^0 \bar{K}^0}$

Time dependent CP asymmetry in  $B_d \rightarrow K^0 \bar{K}^0$



**SM prediction** [A.K.Giri and R.Mohanta, JHEP 0411 (2004) 084]

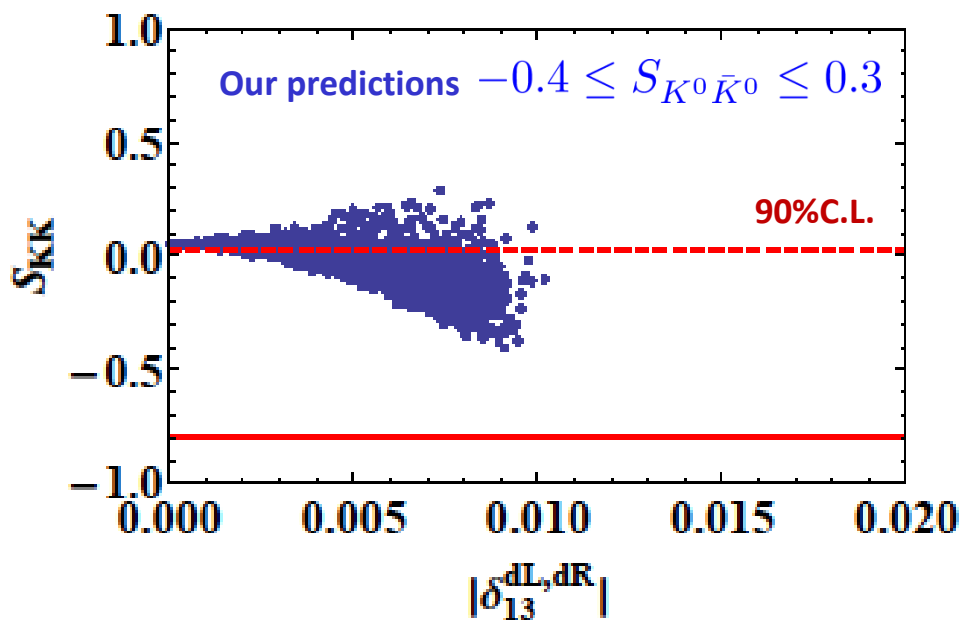
$$S_{K^0 \bar{K}^0}(\text{SM}) \simeq 0.06$$

with pQCD

**Experimental result** [PDG 2012]

$$S_{K^0 \bar{K}^0}(\text{exp}) = -0.8 \pm 0.5$$

**Our prediction**



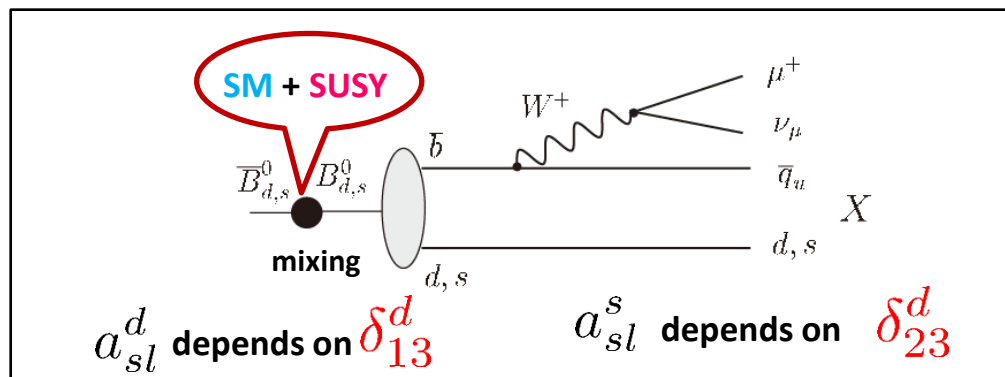
# Predictions

$$a_{sl}^d, a_{sl}^s$$

CP asymmetry in the semileptonic decay  $\bar{B}_q^0 \Rightarrow B_q^0 \rightarrow \mu^+ X$

$$a_{sl}^q \equiv \frac{\Gamma(\bar{B}_q^0 \rightarrow \mu^+ X) - \Gamma(B_q^0 \rightarrow \mu^- X)}{\Gamma(\bar{B}_q^0 \rightarrow \mu^+ X) + \Gamma(B_q^0 \rightarrow \mu^- X)}$$

$$= \frac{1 - \left| \frac{p}{q} \right|^4}{1 + \left| \frac{p}{q} \right|^4} = -\text{Im} \left( \frac{\Gamma_{12}}{M_{12}} \right)$$



## SM predictions

[A.Lenz and U.Nierste, arXiv:1102.4274 [hep-ph]]

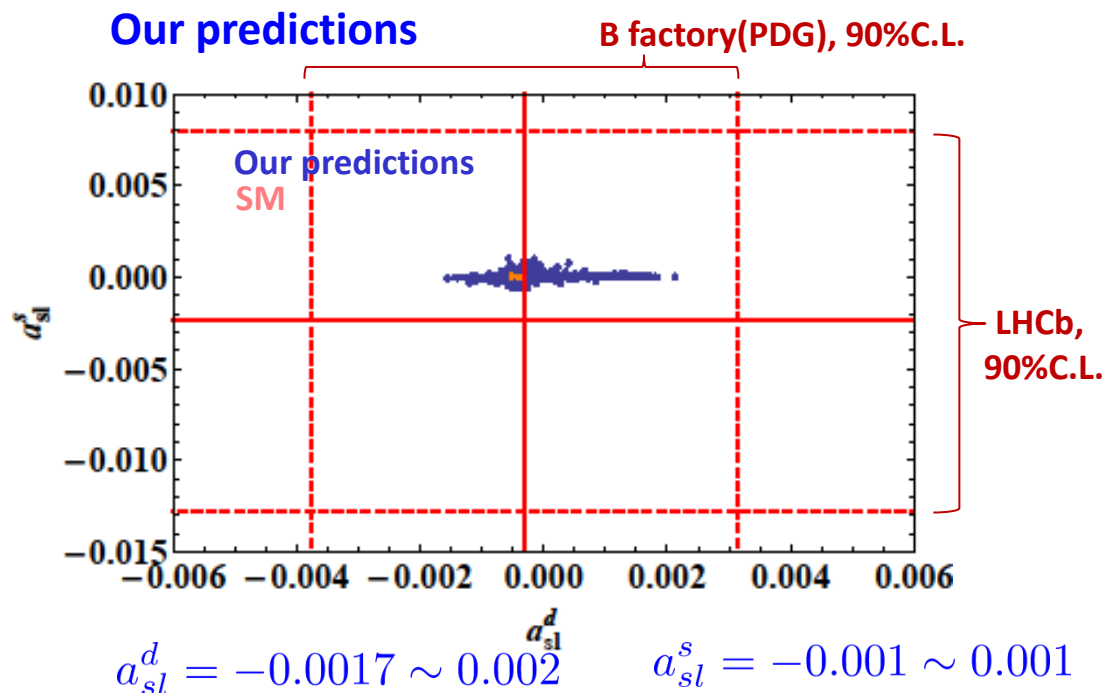
$$a_{sl}^{s, \text{SM}} = (1.9 \pm 0.3) \times 10^{-5}$$

$$a_{sl}^{d, \text{SM}} = -(4.1 \pm 0.6) \times 10^{-4}$$

## Experimental results [PDG 2012]

$$a_{sl}^s = (-0.24 \pm 0.54 \pm 0.33) \times 10^{-2}$$

$$a_{sl}^d = (-0.3 \pm 2.1) \times 10^{-3}$$



# Summary

We have discussed the sensitivity of **the gluino-bottom-quark interaction** to the CP violating phenomena of the  $K$ ,  $B^0$  and  $B_s$  mesons.

We take **the natural SUSY scenario**, which is consistent with the experimental situation of LHC.

$$m_{\tilde{q}_{1,2}} = \mathcal{O}(10) \text{ TeV}, \quad m_{\tilde{b}_1} = 1 \text{ TeV}, \quad m_{\tilde{b}_2} = 1.1 \text{ TeV}, \quad m_{\tilde{g}} = 2 \text{ TeV}$$

The relevant constraints :  $|\epsilon_K|$ ,  $\Delta M_d$ ,  $\Delta M_s$ ,  $S_{J/\psi K_S}$ ,  $S_{J/\psi \phi}$ ,  $\text{BR}(b \rightarrow s\gamma)$ ,  $|d_s^C|$

→ The allowed region of the mixing parameter

$$|\delta_{13}^{dL(dR)}| = 0 \sim 0.01 \quad |\delta_{23}^{dL(dR)}| = 0 \sim 0.04$$

We predict  $S_{\phi K_S}$ ,  $S_{\eta' K^0}$ ,  $S_{\phi\phi}$ ,  $S_{\eta'\phi}$ ,  $S_{K^0 \bar{K}^0}$ ,  $a_{sl}^d$ ,  $a_{sl}^s$

SUSY contributions to  $S_{\phi K_S}$ ,  $S_{\eta' K^0}$  is tiny,

$$\begin{aligned} -0.1 &\leq S_{\phi\phi} \leq 0.2 & -0.1 &\leq S_{\phi\eta'} \leq 0.2 \\ -0.4 &\leq S_{K^0 \bar{K}^0} \leq 0.3 \\ a_{sl}^d &= -0.0017 \sim 0.002 & a_{sl}^s &= -0.001 \sim 0.001 \end{aligned}$$

**Our model will be tested by Belle II as well as LHCb.**



Buckup

## Mixing angle $\theta$

**Mass matrix**

$$\tilde{m}_{\tilde{q}\text{dia}}^2 = \Gamma_G^{(q)} M_{\tilde{q}}^2 \Gamma_G^{(q)\dagger}$$

$$M_{\tilde{q}}^2 = \begin{pmatrix} m_{\tilde{d}_L}^2 & m_b(A_b - \mu \tan\beta) \\ m_b(A_b - \mu \tan\beta) & m_{\tilde{d}_R}^2 \end{pmatrix}$$

**Third family**

$$\tan 2\theta = \frac{2m_b(A_b - \mu \tan\beta)}{m_{\tilde{d}_L}^2 - m_{\tilde{d}_R}^2}$$

$\Delta F=1$

$$H_{eff} = \frac{4G_F}{\sqrt{2}} \left[ \sum_{q'=u,c} V_{q'b} V_{q'q}^* \sum_{i=1,2} C_i O_i^{(q')} - V_{tb} V_{tq}^* \sum_{i=3-6,7\gamma,8G} (C_i O_i + \tilde{C}_i \tilde{O}_i) \right],$$

$$O_1^{(q')} = (\bar{q}_\alpha \gamma_\mu P_L q'_\beta) (\bar{q}'_\beta \gamma^\mu P_L b_\alpha), \quad O_2^{(q')} = (\bar{q}_\alpha \gamma_\mu P_L q'_\alpha) (\bar{q}'_\beta \gamma^\mu P_L b_\beta),$$

$$O_3 = (\bar{q}_\alpha \gamma_\mu P_L b_\alpha) \sum_Q (\bar{Q}_\beta \gamma^\mu P_L Q_\beta), \quad O_4 = (\bar{q}_\alpha \gamma_\mu P_L b_\beta) \sum_Q (\bar{Q}_\beta \gamma^\mu P_L Q_\alpha),$$

$$O_5 = (\bar{q}_\alpha \gamma_\mu P_L b_\alpha) \sum_Q (\bar{Q}_\beta \gamma^\mu P_R Q_\beta), \quad O_6 = (\bar{q}_\alpha \gamma_\mu P_L b_\beta) \sum_Q (\bar{Q}_\beta \gamma^\mu P_R Q_\alpha),$$

$$O_{7\gamma} = \frac{e}{16\pi^2} m_b \bar{q}_\alpha \sigma^{\mu\nu} P_R b_\alpha F_{\mu\nu}, \quad O_{8G} = \frac{g_s}{16\pi^2} m_b \bar{q}_\alpha \sigma^{\mu\nu} P_R T_{\alpha\beta}^a b_\beta G_{\mu\nu}^a,$$

# ΔF=1

Most dominant term comes from  $C_{8G}^{\tilde{g}}$ .

$$A^{SUSY} (\bar{B}_d \rightarrow \phi K_s) \propto C_{8G}^{\tilde{g}}(m_b) + \tilde{C}_{8G}^{\tilde{g}}(m_b)$$

$$\tilde{C} : C(L \Leftrightarrow R)$$

$$C_{8G}^{\tilde{g}}(m_{\tilde{g}}) = \frac{8}{3} \frac{\sqrt{2}\alpha_s\pi}{2G_F V_{tb} V_{tq}^*} \left[ \frac{(\Gamma_{GL}^{(d)})_{k3}^*}{m_{\tilde{d}_3}^2} \left\{ (\Gamma_{GL}^{(d)})_{33} \left( -\frac{9}{8} F_1(x_{\tilde{g}}^3) - \frac{1}{8} F_2(x_{\tilde{g}}^3) \right) \right. \right. \\ \left. \left. + \frac{m_{\tilde{g}}}{m_b} (\Gamma_{GR}^{(d)})_{33} \left( -\frac{9}{8} F_3(x_{\tilde{g}}^3) - \frac{1}{8} F_4(x_{\tilde{g}}^3) \right) \right\} \right. \\ \left. + \frac{(\Gamma_{GL}^{(d)})_{k6}^*}{m_{\tilde{d}_6}^2} \left\{ (\Gamma_{GL}^{(d)})_{36} \left( -\frac{9}{8} F_1(x_{\tilde{g}}^6) - \frac{1}{8} F_2(x_{\tilde{g}}^6) \right) \right. \right. \\ \left. \left. + \frac{m_{\tilde{g}}}{m_b} (\Gamma_{GR}^{(d)})_{36} \left( -\frac{9}{8} F_3(x_{\tilde{g}}^6) - \frac{1}{8} F_4(x_{\tilde{g}}^6) \right) \right\} \right], \quad x = \frac{m_{\tilde{g}}^2}{m_{\tilde{d}_i}^2}$$

$\delta_{23}^{dL} s_\theta c_\theta e^{-i\phi}$

$-\delta_{23}^{dL} s_\theta c_\theta e^{-i\phi}$

If two masses  $m_{\tilde{b}_1}, m_{\tilde{b}_2}$  degenerate, we have complete cancellation.

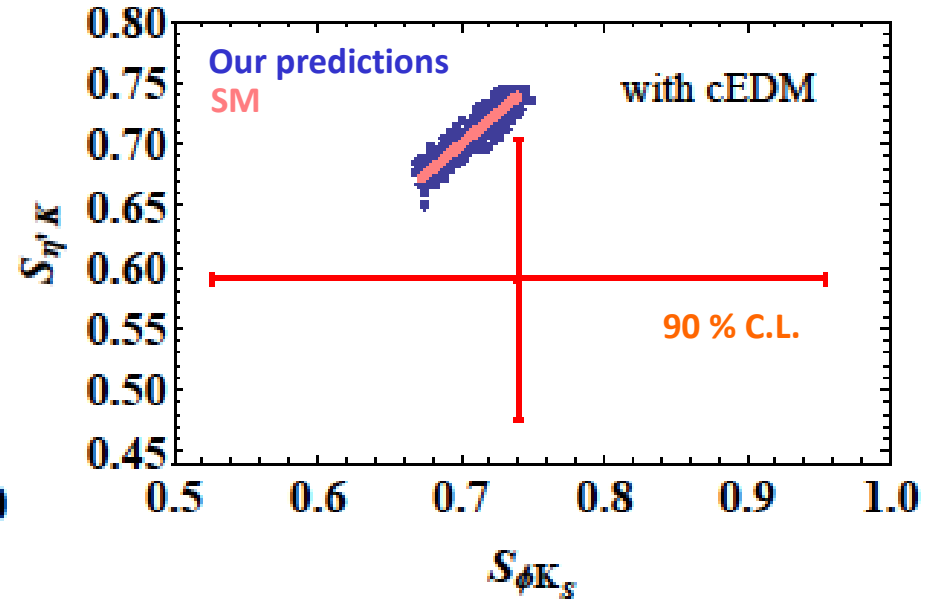
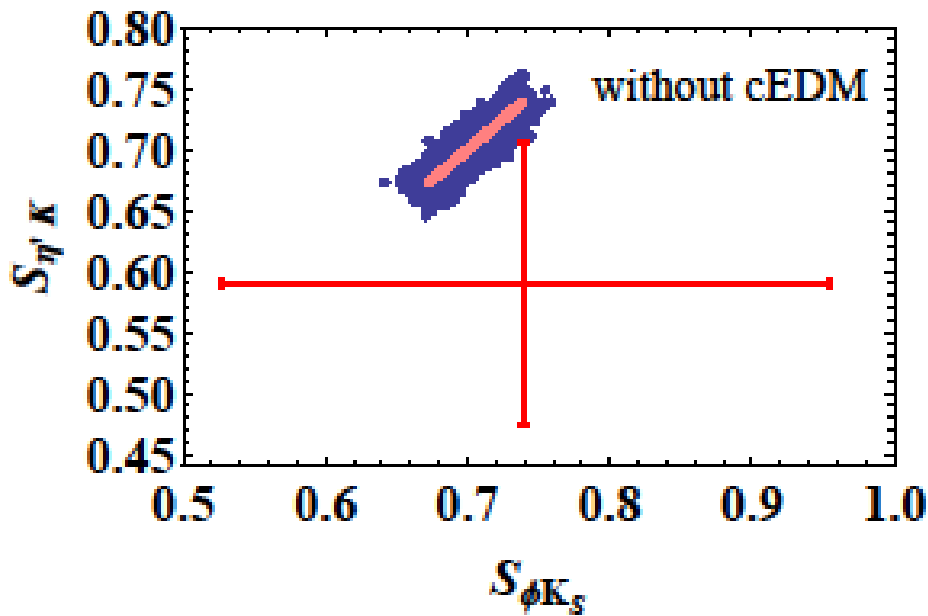
$C_{8G}^{\tilde{g}}$  is reduced by this cancellation.

## Factorization relation

$$\langle O_3 \rangle = \langle O_4 \rangle = \left(1 + \frac{1}{N_c}\right) \langle O_5 \rangle, \quad \langle O_6 \rangle = \frac{1}{N_c} \langle O_5 \rangle,$$

$$\langle O_{8G} \rangle = \frac{\alpha_s(m_b)}{8\pi} \left( -\frac{2m_b}{\sqrt{\langle q^2 \rangle}} \right) \left( \langle O_4 \rangle + \langle O_6 \rangle - \frac{1}{N_c} (\langle O_3 \rangle + \langle O_5 \rangle) \right),$$

# The effect of cEDM of the strange quark



## Input parameters

$$\alpha_s(M_Z) = 0.1184 \text{ [34]}$$

$$m_c(m_c) = 1.275 \text{ GeV [34]}$$

$$m_t(m_c) = 1.275 \text{ GeV } (\overline{MS}) \text{ [34]}$$

$$M_{B_s} = 5.36677(24) \text{ GeV [34]}$$

$$\Delta M_s = (116.942 \pm 0.1564) \times 10^{-13} \text{ GeV [7]}$$

$$\Delta M_d = (3.337 \pm 0.033) \times 10^{-13} \text{ GeV [34]}$$

$$f_{B_s} = (233 \pm 10) \text{ MeV [47]}$$

$$f_{B_s}/f_{B^0} = 1.200 \pm 0.02 \text{ [47]}$$

$$\xi_s = 1.21(6) \text{ [26]}$$

$$\lambda = 0.2255(7) \text{ [34]}$$

$$|V_{cb}| = (4.12 \pm 0.11) \times 10^{-2} \text{ [47]}$$

$$\eta_{cc} = 1.43(23) \text{ [26]}$$

$$\eta_{ct} = 0.47(4) \text{ [26]}$$

$$\eta_{tt} = 0.5765(65) \text{ [26]}$$

$$f_K = (156.1 \pm 1.1) \text{ MeV [34]}$$

$$\kappa_\epsilon = 0.92(2) \text{ [26]}$$

# The unitarity triangle

